

CHAPTER 29: Magnetic Fields and Sources



*photo courtesy of
Mr. White*

Relativity and Magnetism

--*Newton's Classical Mechanics* demands that the **speed of light** depend upon the **relative motion** between the *frame of reference* of the light source and the observer's frame (think of about a passing car's speed on the freeway—how fast it passes you depends upon how fast *you* are going).

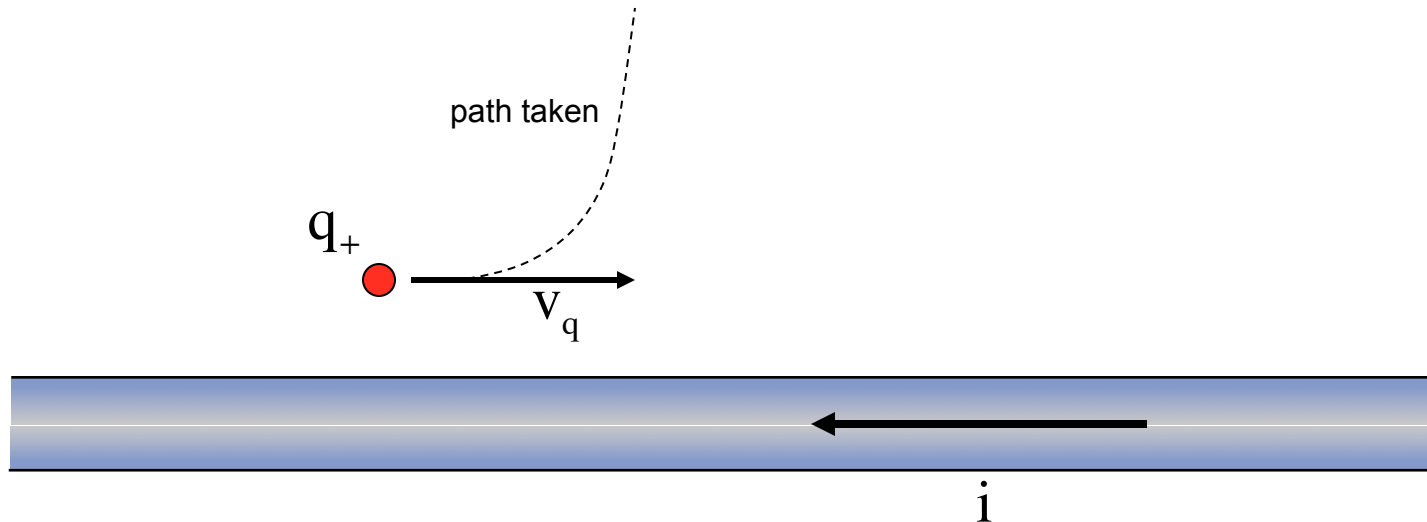
--*Yet because* **alternating magnetic fields (B-flds) induce electric fields (E-flds)**, and **alternating E-flds induce B-flds**, the only way an alternating **E-fld can coupled** with an alternating **B-fld** to **produce an electromagnetic wave** (i.e., a wave in which the alternating E-fld feeds the B-fld and the alternating B-fld feeds the E-fld), is if, according to Maxwell's equations, the wave's velocity is 3×10^8 m/s, *the speed of light*.

--*You can't have it* both ways. The *speed of light* is **either frame-of-reference dependent** alla Newton, **or** it's a **fixed value independent of frame-of-reference** alla Maxwell.

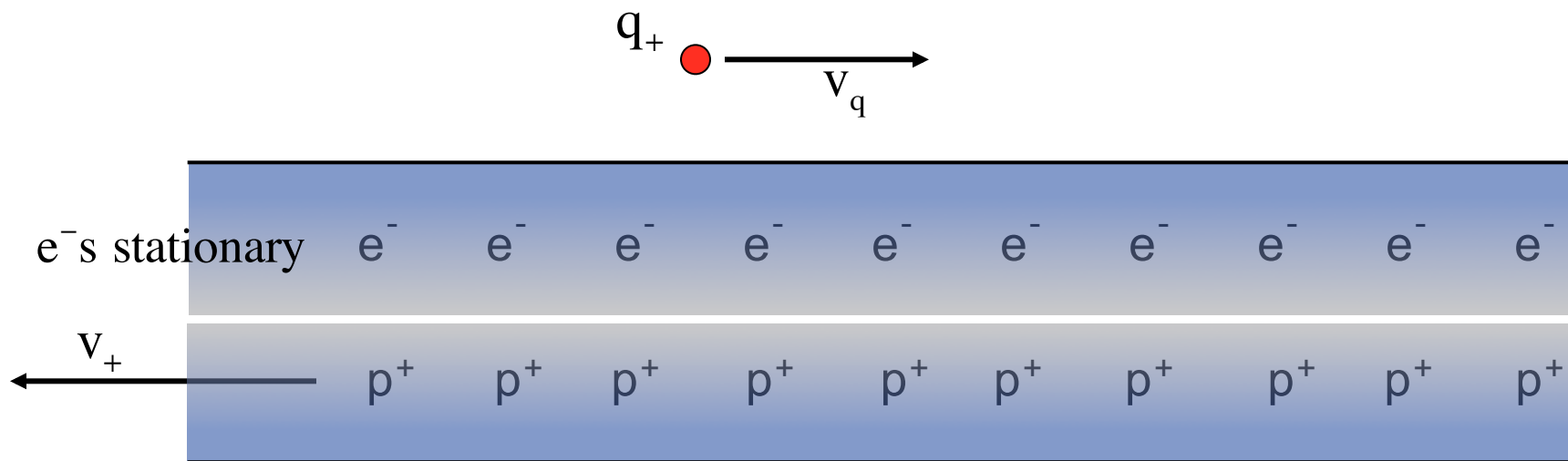
--*This conundrum* is **what motivated Einstein** to develop his **Theory of Special Relativity**, a form of Mechanics that **assume that the speed of light is the same in all frames of reference**.

--One of the stranger characteristics of *Special Relativity* is that if you have an object that is approaching at relativistic speeds, it will *length contract*. And, in fact, this length contraction phenomenon will occur even at classical speeds (though observing it at classical speeds is difficult).

--With this in mind, consider a wire with conventional current flowing through it (i.e., assume positive charge flow). If you fire a positive charge opposite the direction of current flow, an interesting thing is observed. The charge will feel a force that motivates it to veer away from the wire. So what's going on?

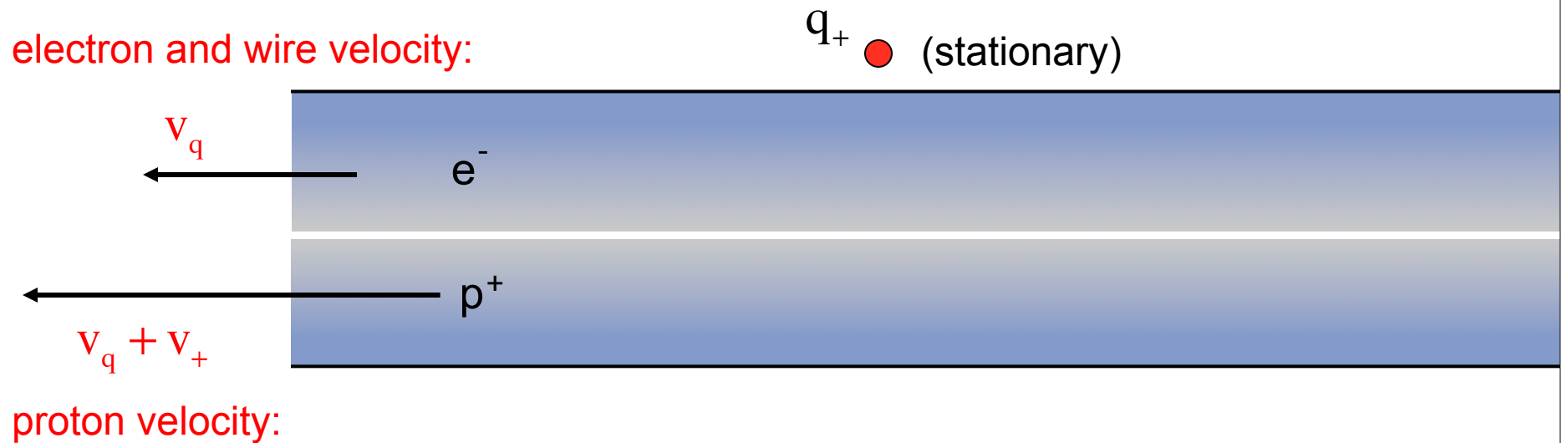


As *positive charge* carriers move onto the wire (with electrons assumed stationary), an equal number of positive charge carriers move off the wire. That means the number of positive and negative charge stays even throughout time, the wire stays *electrically neutral*, and there's *no good reason for the moving test charge to feel a force*. But it does. This led early theorists to conclude that there must be a force, a *magnetic force*, affecting the moving charge . . . except *that isn't what's really happening* here. To see this, we have to look at the situation through the perspective of the moving charge q_+ .

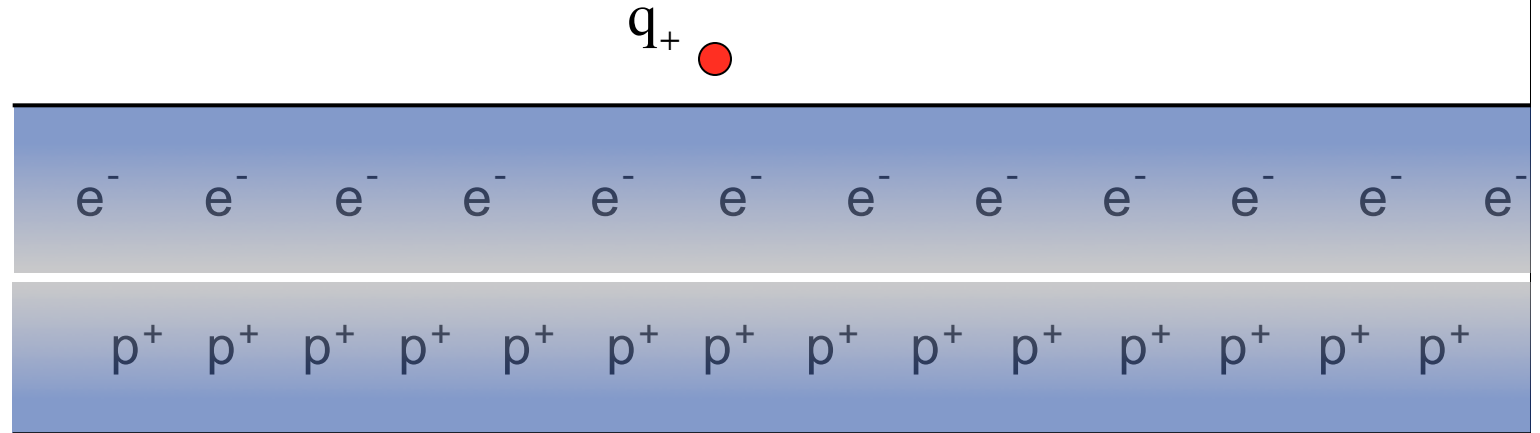


assume velocity of positive charge is v_+

In the frame of reference attached to q_+ , q_+ is **not moving**. What's more, as far as q_+ is concerned, **both the electrons and the wire are seen to be moving to the left with velocity v_q** , and the **protons are seen to be moving to the left with velocity $v_q + v_+$** . The sketch below show all of this.

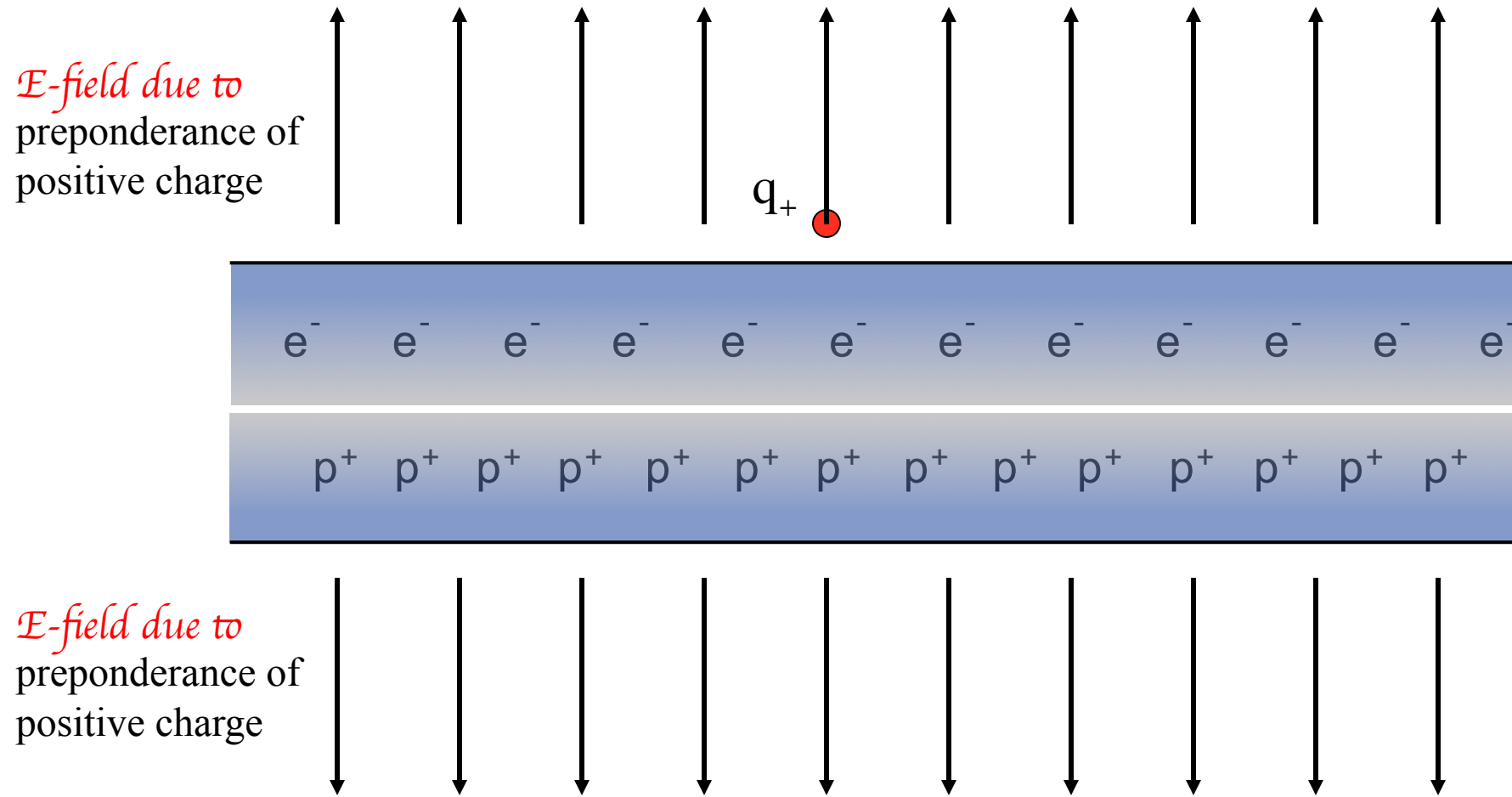


What's important to notice here is that *because the protons are moving faster than the electrons*, they will *length contract more* than will the electrons. When they do so, from the moving charge's perspective, there appears to be *more protons on the wire than electrons*.



Protons length contract more than electrons. More protons means the *wire looks electrically positive* and an *electric field is set up pointing outward* from the wire.

With more protons apparently on the wire, there is an electric field generated emanating outward from the wire. It is that electric field that the q_+ responds to.



In short, what was described by early researchers as a **MAGNETIC EFFECT** was (and is) really a **RELATIVISTIC EFFECT**.

Still, the **Classical Theory of Magnetism** is a good theory in the everyday world (just as is the case with Newtonian Mechanics), so that's what you will be spending the next several weeks learning.

General Information

Electric Fields

--*electric fields* (abbreviated as *E-flds*), with units of *newtons per coulomb* or *volt per meter*, are **modified force fields** (release a charge in an E-fld and it will accelerate);

--*electric fields* are generated with the **presence of charge**;

--an *electric field's direction* is defined as the **direction a positive charge will accelerate if released in the field**;

--*electric field lines*:

--go from **positive to negative charge**;

--*identify* the E-fld's **direction** in a region;

--*are closer together* where E-flds are more intense;

Magnetic Fields

--*magnetic fields* (abbreviated as *B-flds*), with units of *teslas* in the *MKS system*, are **NOT modified force fields** (release a charge in a B-fld and it will just sit there);

--*magnetic forces do exist* when a **charge moves through a B-fld**—they are **centripetal** and are governed by the relationship: $\vec{F} = q\vec{v} \times \vec{B}$

--*B-fields* are generated by **charge in motion**;

--a *B-field's direction* is defined as the **direction a compass points when placed in the field**;



--*magnetic field lines*:

--go from **north to south pole, or circle around current carrying wire**;

--*identify* the B-fld's **direction** in a region;

--*are closer together* where B-flds are more intense;

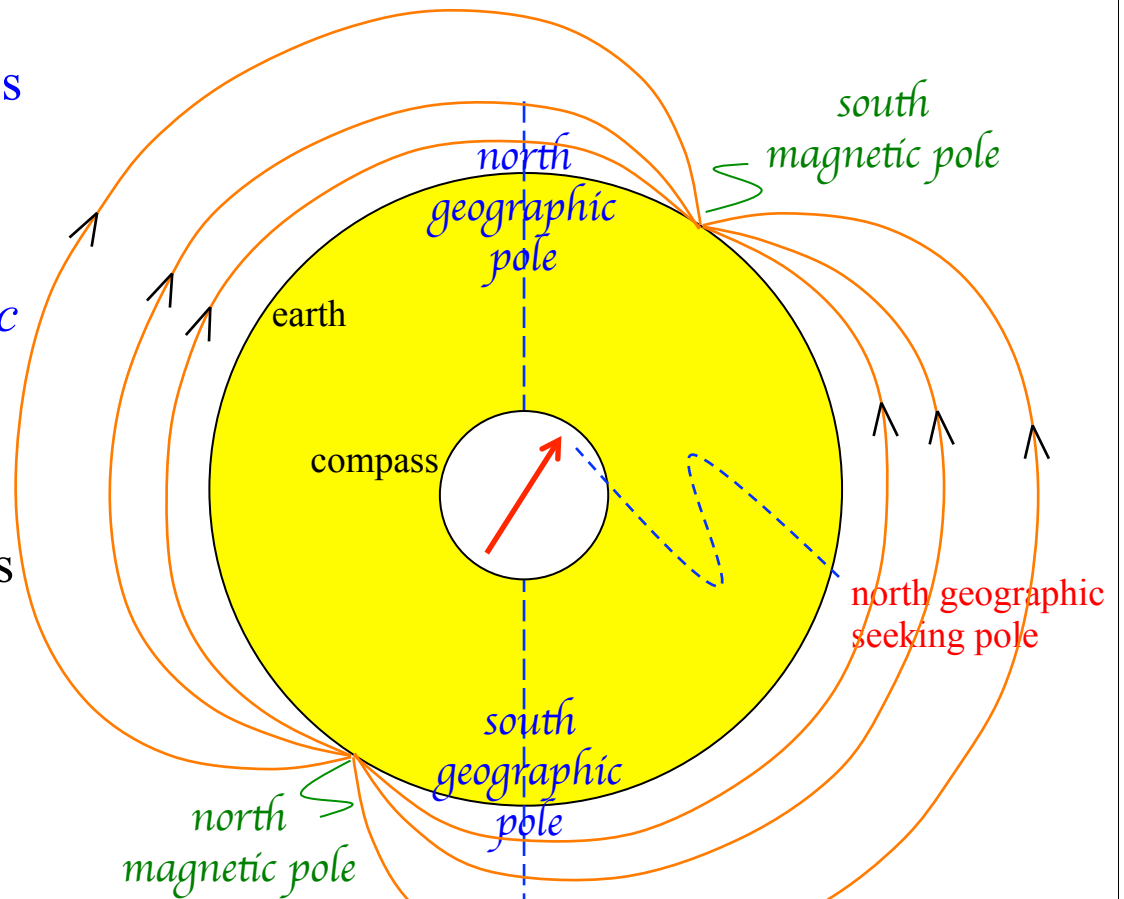
Terminology and the Compass

Place a compass in the earth's magnetic field.

The compass end that originally pointed toward the north geographic hemisphere was called the north geographic seeking magnetic pole.

With time, the word geographic was dropped leaving the north seeking magnetic pole, and with even more time, the seeking and magnetic was dropped leaving us with north pole.

Problem is, north poles are attracted to south poles, which means that given the definition, there must be a south magnetic pole in the north geographic hemisphere. Not very esthetically pleasing, but that's life (and it'll switch directions in another 200,000 or so years due to slow oscillatory patterns in the earth's magma).



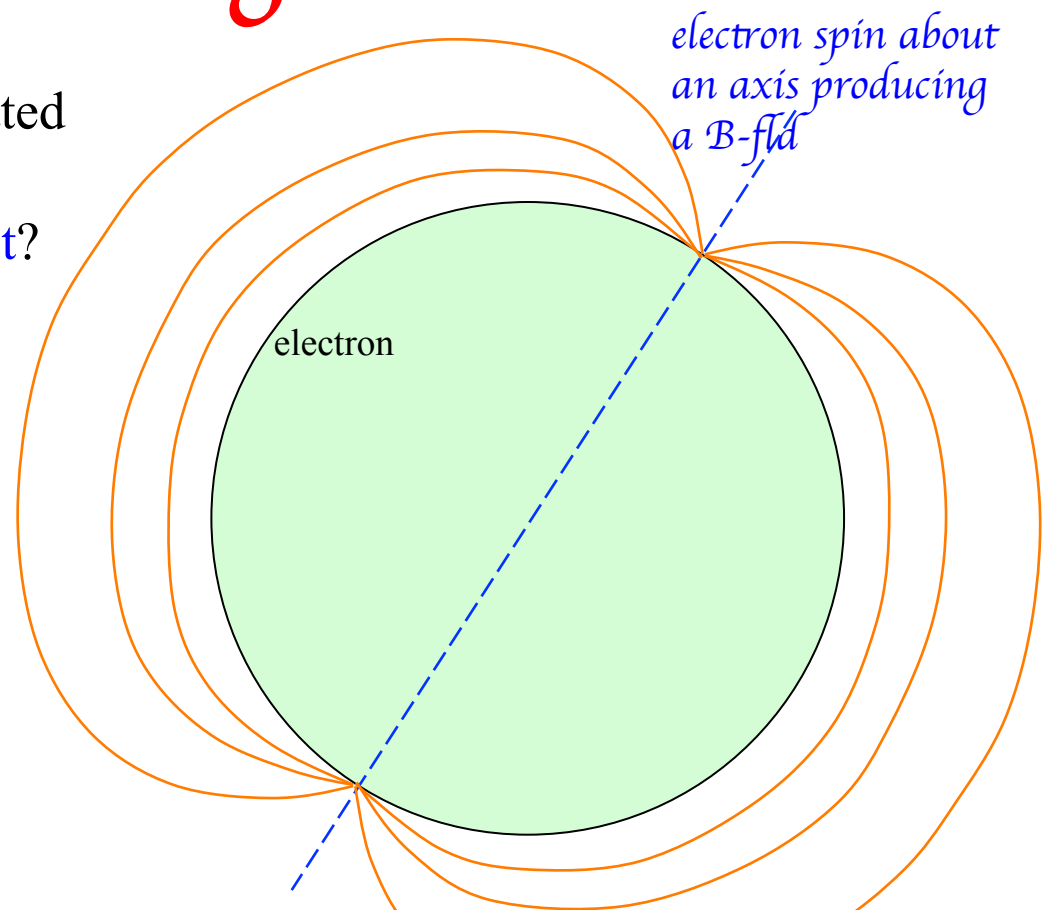
Bar Magnets

If magnetic fields are generated by charge in motion, where is the motion associated with a bar magnet?

An atom's spin quantum number highlights the fact that electrons spin up or down, depending.

In most atoms, approximately the same number of electrons spin in one direction as the other, but in certain atoms (the ferromagnetic ones), they spin considerably more in one direction than the other.

The most prominent example of a ferromagnetic material is iron, with six more electrons spinning in one direction than the other. As such, EACH IRON ATOM IS A MINI-MAGNET UNTO ITSELF.

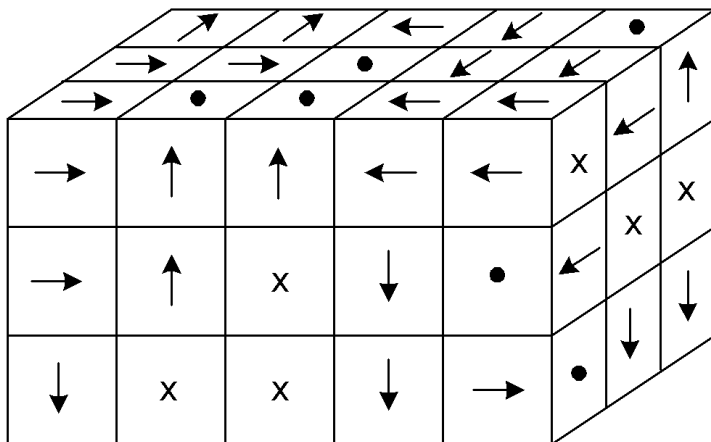


So how can a piece of iron *be magnetic under some conditions* and not under other conditions (there are, after all, iron nails that do not exhibit magnetic characteristics at all). (The explanation is called *Ampere's Theory of Magnetism*.)

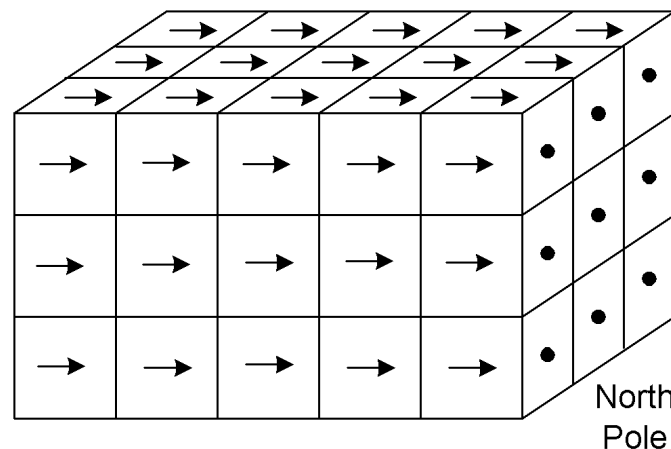
Enter the magnetic domain. What happens is this: Atoms within irregular, microscopic regions, called domains, align themselves so that all of their magnetic fields are in the same direction.

When the domains are themselves aligned, the material acts like a magnet. When the domains are NOT aligned, the *net magnetic effect is lost* (in some cases, all that is needed to de-magnetize a magnetic is to have thermal agitation shake the domains out of alignment). In any case, the first sketch below is without alignment, the second with alignment.

domains unaligned (not realistic rendering)



domains aligned



sketches courtesy of Mr. White

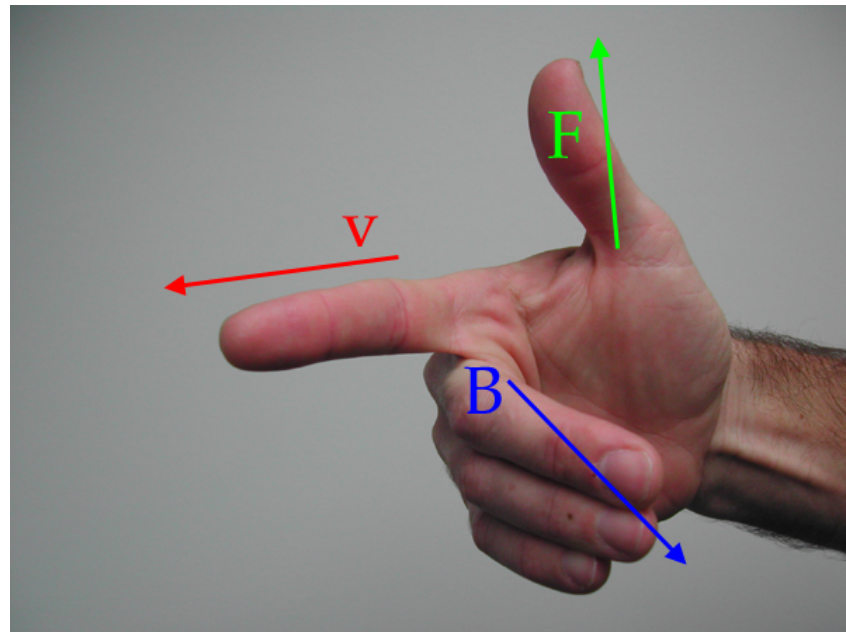
Magnetic Force

When charge moves through a magnetic field, it may or may not feel a force, depending upon its motion. If present, that force will be:

$$\vec{F} = q\vec{v} \times \vec{B}$$

The magnitude is $|\vec{F}| = q|\vec{v}||\vec{B}|\sin\theta$, where q is the size of the charge, $|\vec{v}|$ is the magnitude of the velocity vector, $|\vec{B}|$ is the magnitude of the magnetic field and θ is the angle between the line of the two vectors.

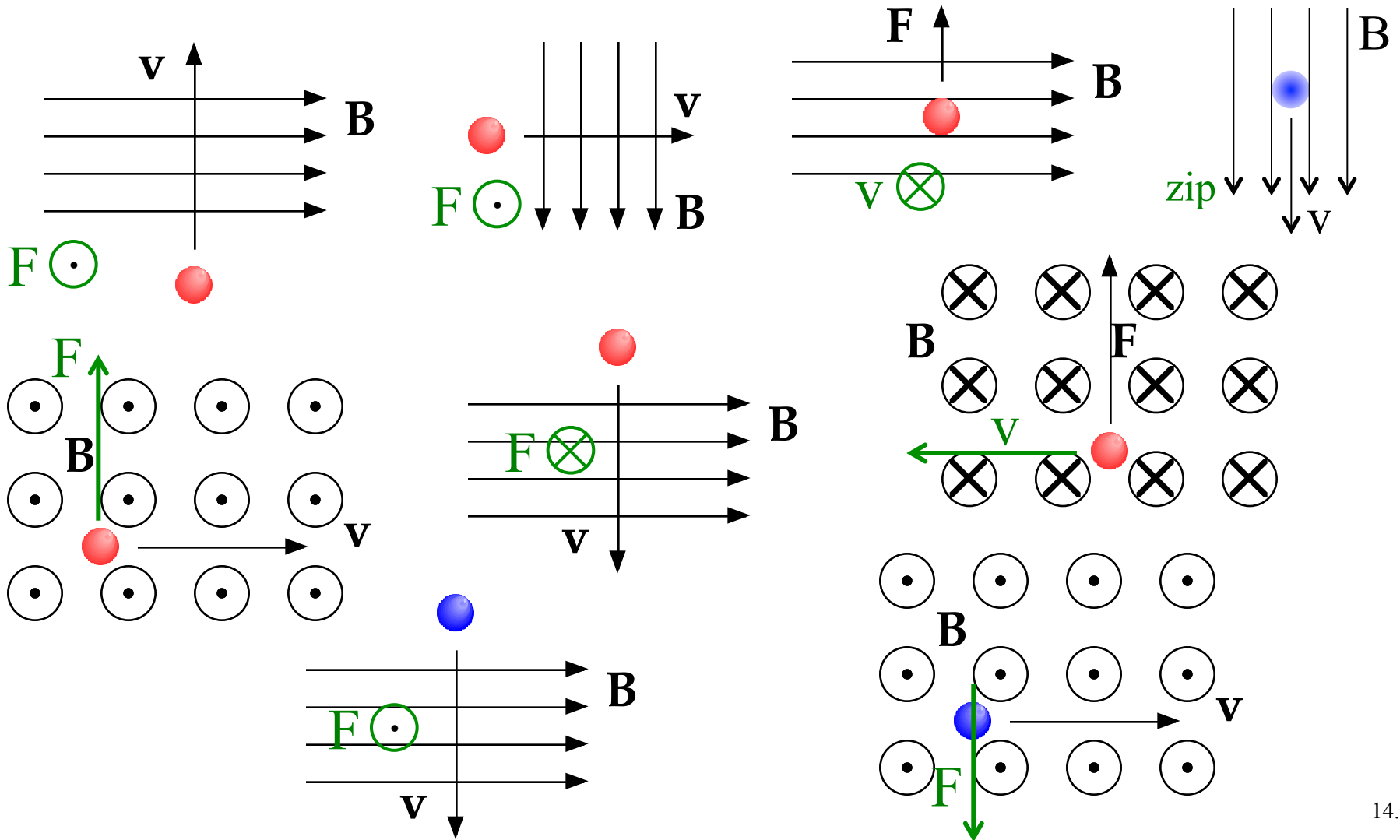
The direction is determined using the *right-hand rule*.



sketch courtesy of
Mr. White

Example 1: (courtesy of Mr. White)

Identify the missing vector in these diagrams. Assume red signifies negative charges, blue positive charges. The responses are in green



Fine Print for $\vec{F} = q\vec{v} \times \vec{B}$

The *right-hand-rule* determine the *direction of force* for a *positive charge*.

Magnetic forces are *centripetal forces* acting *perpendicular* to magnetic fields (whereas electric forces act along electric fields). That means magnetic forces **DO NO WORK** on charges that feel their effect.

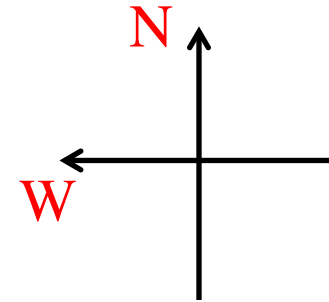
Magnetic forces are experience *only by charges in motion*.

Example 2: A proton moving upward with speed 5×10^6 m/s in a magnetic field feels a force of 8×10^{-14} N to the west. When moving horizontally to the north, it feels no force. Find the magnitude and direction of the magnetic field in this region. (courtesy of Mr. White)

Because no force is felt when the proton is moving northward, the magnetic field must either be toward the north or the south.

Reverse engineering the right-hand-rule suggests if the field is to the north*, the velocity vector would have to be out of the page to generate a force to the west, with a magnitude of:

$$\begin{aligned} F &= qvB \sin 90^\circ \\ \Rightarrow B &= \frac{F}{qv} \\ &= \frac{(8 \times 10^{-14} \text{ N})}{(1.6 \times 10^{-19} \text{ C})(5 \times 10^6 \text{ m/s})} \\ \Rightarrow &= 10^{-1} \text{ T} \end{aligned}$$



*Note that if the field was to the south, the velocity vector would be into the page.

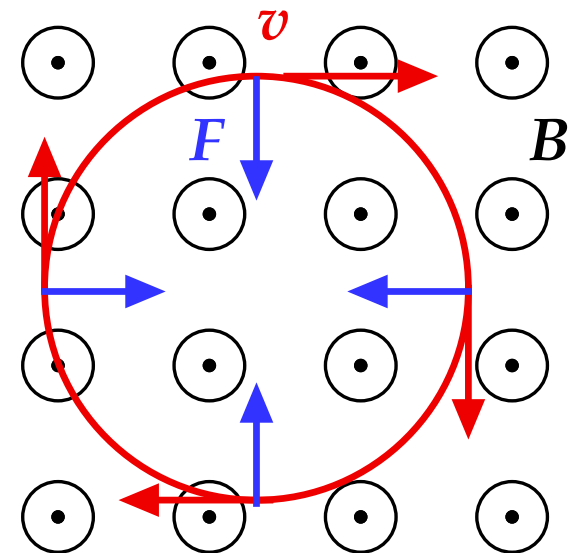
Example 3: A charge q of mass m is moving with constant velocity v at right angles to a magnetic field B . (idea courtesy of Mr. White)

a.) What kind of motion will it execute?

Because magnetic forces are centripetal, the mass will follow a circular path.

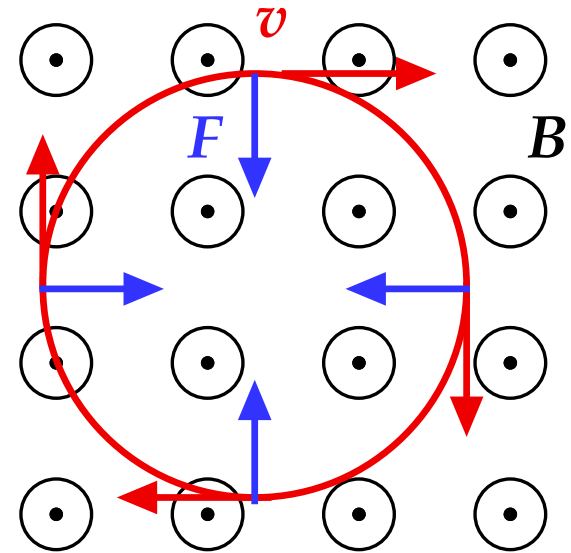
b.) What is the radius of the motion's path?

$$\begin{aligned}\sum F_{\text{cent}} : \\ q\vec{v} \times \vec{B} &= m\vec{a}_{\text{cent}} \\ \Rightarrow qvB \sin 90^\circ &= m \frac{v^2}{R} \\ \Rightarrow R &= \frac{mv}{qB}\end{aligned}$$



c.) What is the period of the motion?

$$\begin{aligned} T &= \frac{1}{\nu} \\ &= \frac{1}{\left(\frac{\omega}{2\pi}\right)} = \frac{1}{\left(\frac{\left(\frac{v}{R}\right)}{2\pi}\right)} \\ &= \frac{2\pi}{\left(\frac{v}{R}\right)} \\ &= \frac{2\pi}{\cancel{v} / \left(\frac{mv}{qB}\right)} \\ \Rightarrow T &= 2\pi \frac{m}{qB} \end{aligned}$$

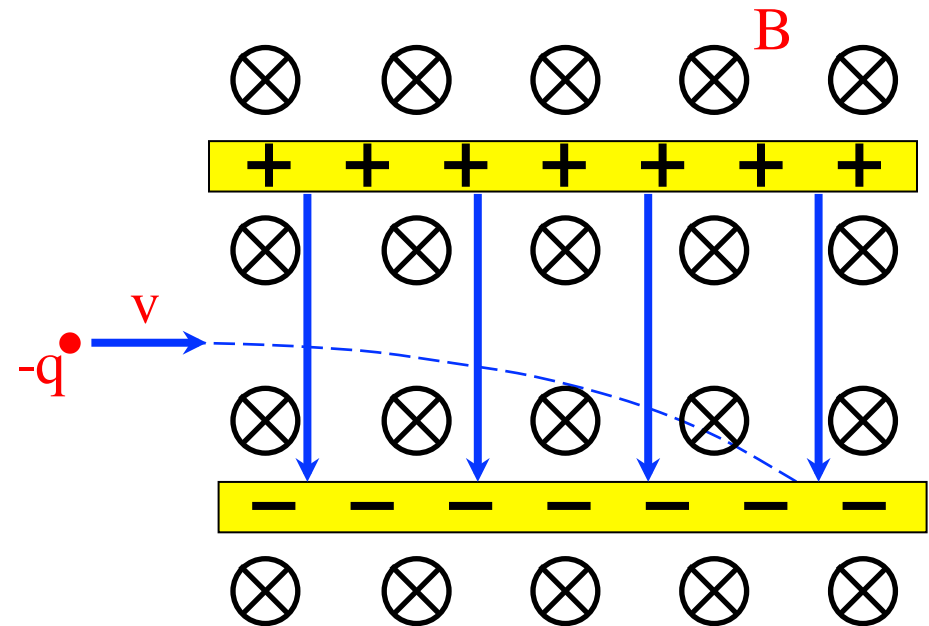


Lorentz Relationship

When both magnetic and electric forces act, the relationship looks like:

$$\begin{aligned}\vec{F}_{\text{net}} &= \mathbf{F}_E + \mathbf{F}_B \\ &= q\vec{E} + q\vec{v}\times\vec{B}\end{aligned}$$

Example 4 (the velocity trap): A velocity trap (or velocity selector) is an electric field and magnetic field at right angles to one another that selectively allow charged particles with *one velocity only* to proceed down its axis. It is made up of parallel plates bathed in a B-fld (see sketch).



a.) Assuming the *B-fld* is into the page as shown, what path will a **negative charge** take if **nothing** additional is added to the system (i.e., no *E-fld* is present).

According to the right-hand-rule, a positive charge would feel a force *upward*, so a **negative charge** will feel a force *downward*.

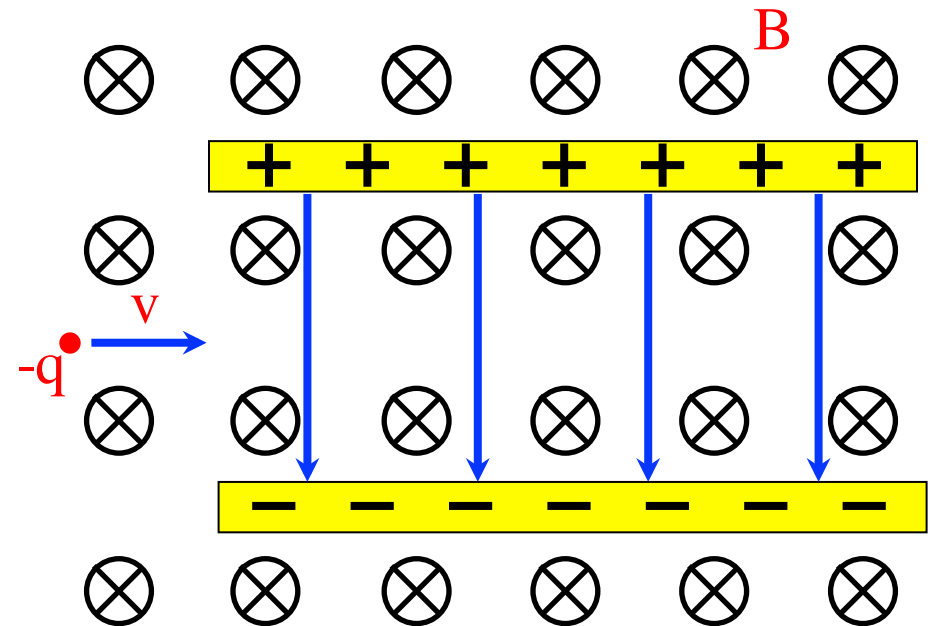
b.) Insert the *E-fld* required to make the charge travel along a straight line.

The force exerted by the *B-fld* is *downward*, so the force exerted by the *E-fld* must be *upward*. To create an *upward* force on a *negative charge*, you need an *E-fld* that is *downward*.

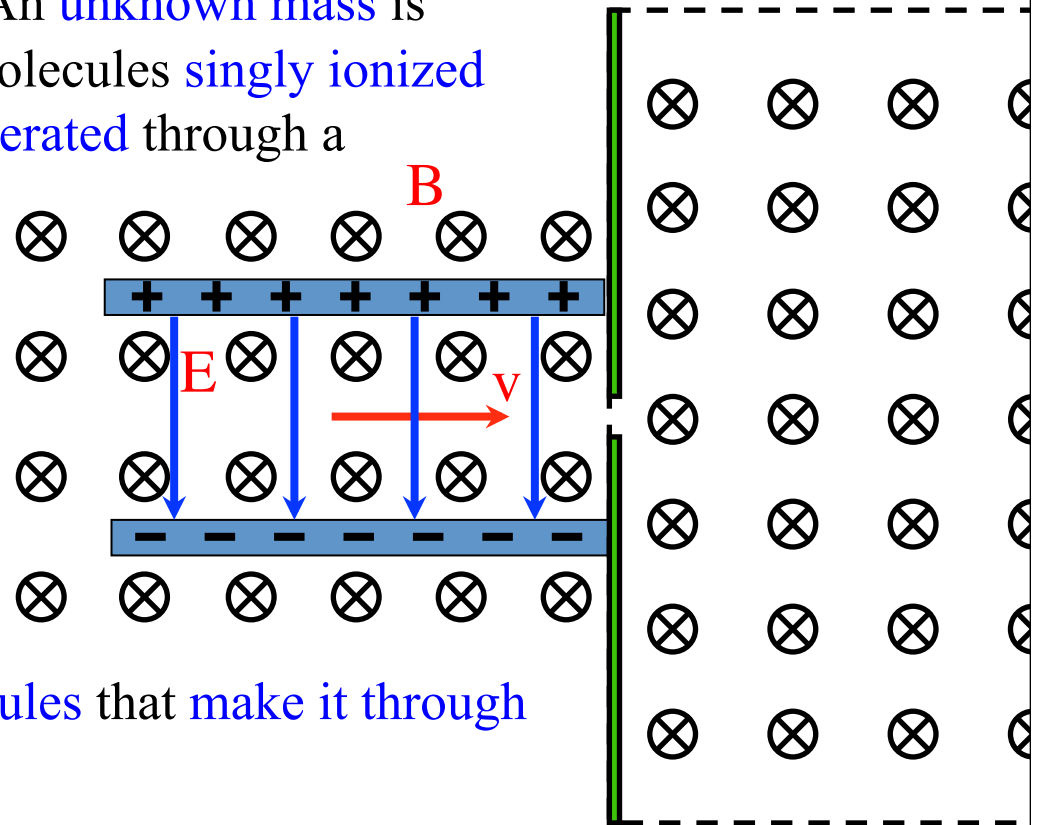
c.) *What velocity* will make it through the trap?

This is a Lorentz relationship problem. The net force must be zero for this to work, which means:

$$qE = qvB$$
$$\Rightarrow v = \frac{E}{B}$$



Example 5 (mass spectrometer): An unknown mass is volatilized (made into a gas), had its molecules singly ionized (had one electron stripped away), accelerated through a potential difference to give them velocity, and sent through a velocity trap made up of a $95,000 \text{ V/m}$ E -fld and a $.93 \text{ teslas}$ B -fld. The molecules that make it through the trap move into a region in which there is only the B -fld.



a.) What is the velocity of the molecules that make it through the trap?

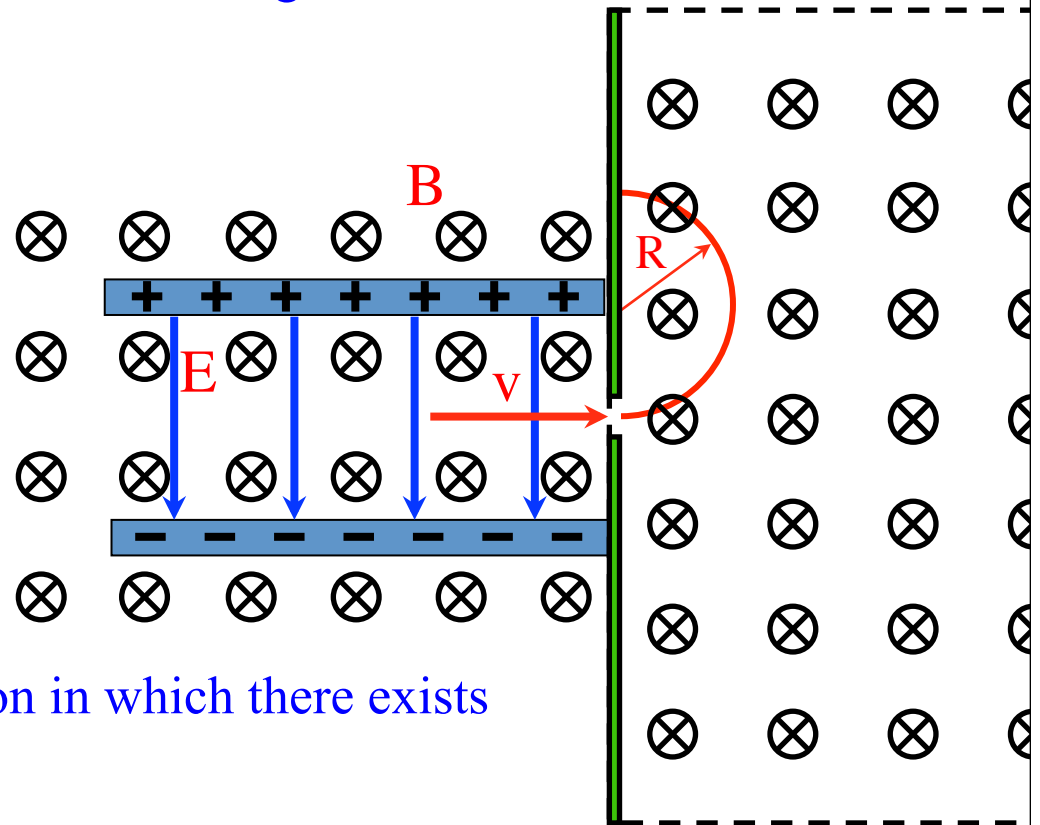
$$qE = qvB$$

$$\Rightarrow v = \frac{E}{B} = \frac{9.5 \times 10^4 \text{ V/m}}{.93 \text{ T}} = 1.02 \times 10^5 \text{ m/s}$$

b.) Draw in the path of the molecules in the far right chamber.

They will circle upward.

c.) If the radius of the arc is observed to be .0667 meters, what was the mass of the particle (that's what these devices are designed to do—determine the mass of an unknown material from which the identification of the material can be had). In the region in which there exists only a B-fld:



$$qvB = m \frac{v^2}{R}$$

$$\Rightarrow m = \frac{qBR}{v} = \frac{(1.6 \times 10^{-19} \text{ C})(.93 \text{ T})(6.67 \times 10^{-2} \text{ m})}{(1.02 \times 10^5 \text{ m/s})}$$

$$= 9.7 \times 10^{-26} \text{ kg}$$

(This is the molecular weight of table salt.)

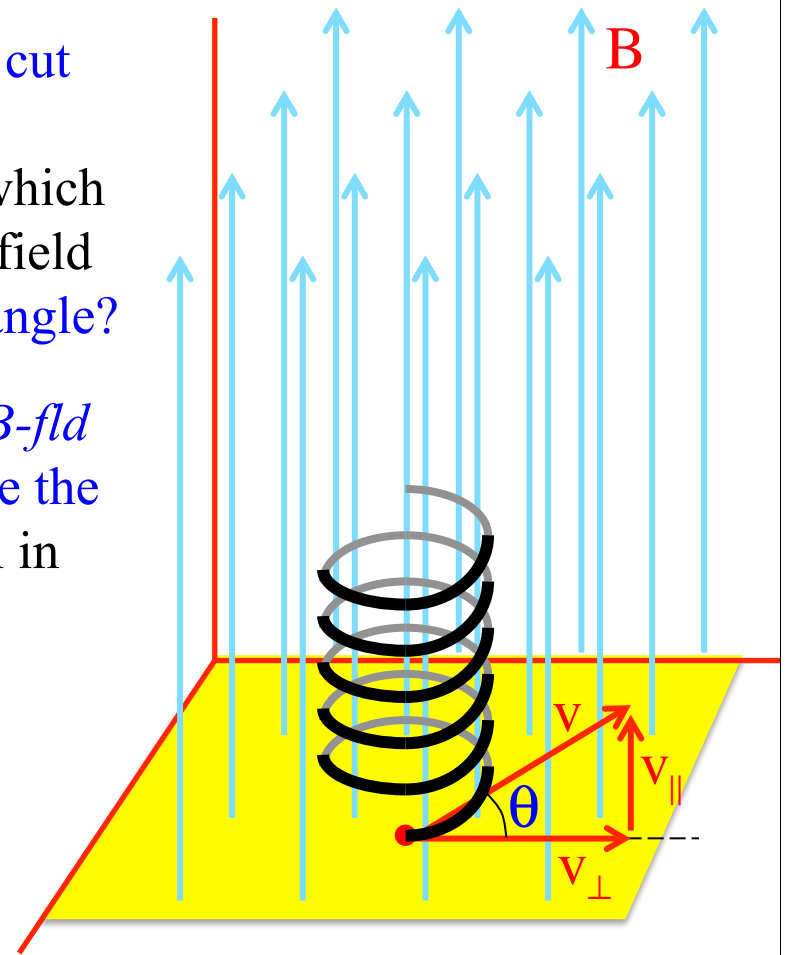
Point of Order

Magnetic forces exist on charges moving through magnetic fields **ONLY** when the charges cut across magnetic field lines. Until now, the only examples we've viewed have been situations in which the charges have cut across at right-angles to the field lines. What happens when they cut across at an angle?

--*The velocity component* perpendicular to the *B-field* will generate a magnetic force that will motivate the charge to circle (magnetic forces *are* centripetal in nature);

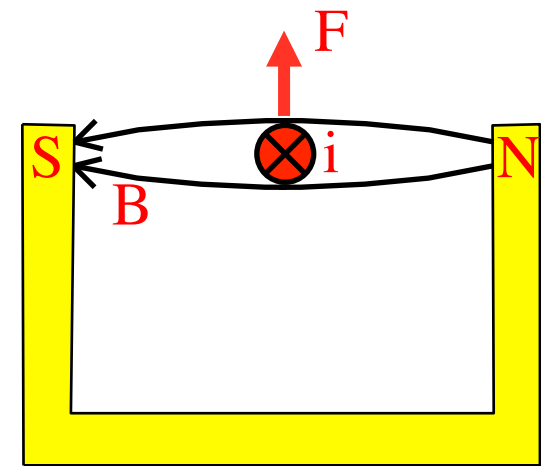
--*The parallel component* will simply provide momentum for the charge to continue to move in that *parallel* direction.

--*The net effect* is that the charge will helix along the *B-field* lines.



B-flds and Current Carrying Wires

Consider a current moving into the page that is positioned between the poles of a horseshoe magnet. What effect will the wire feel due to being in the B-fld.

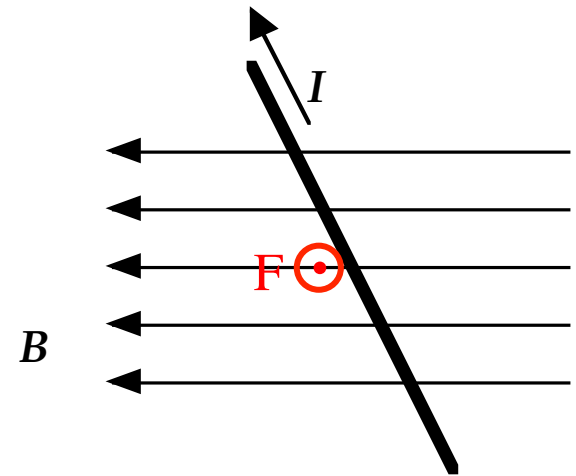


--According to the right-hand-rule, conventional current is the motion of positive charge. Positive charge into the page in a magnetic field to the left produces a magnetic force upward.

--Mathematically: Noticing that a charge's velocity v is the distance it travels L per unit time, and current (q/t) is the charge q that passes by per unit time, we can write:

$$\begin{aligned}\vec{F} &= q\vec{v}\times\vec{B} \\ \Rightarrow \vec{F}_{\text{wire}} &= q\left(\frac{\vec{L}}{t}\right)\times\vec{B} \\ &= i\vec{L}\times\vec{B}\end{aligned}$$

Example 6: (courtesy of Mr. White): A 12 cm length of wire carrying a 30-Amp current is placed in a magnet field at an angle of 60° relative to the field's direction (as shown). If the field is 0.90 T, what force does the wire feel, and in which direction?



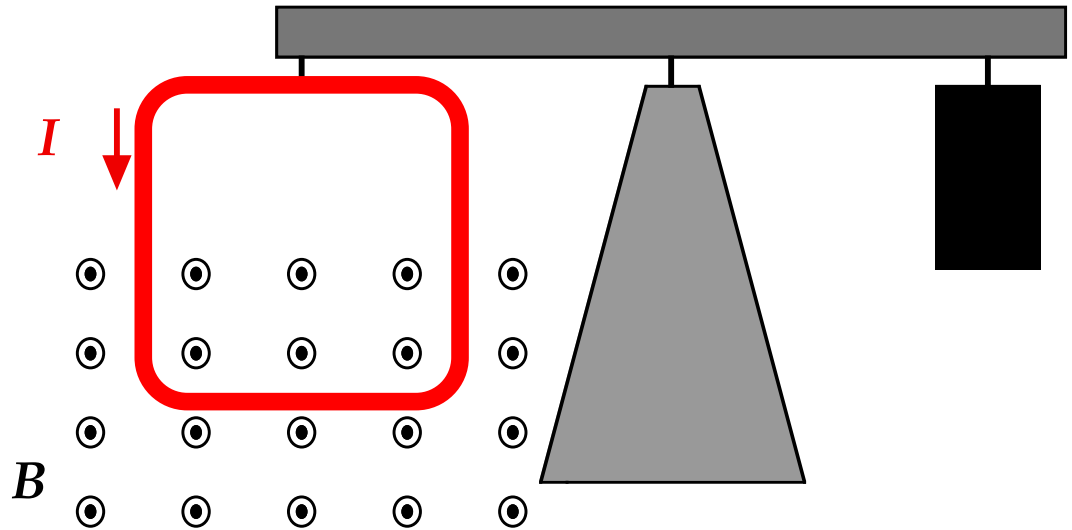
a.) On the sketch, show the direction of the force on the wire.

According to the r.h.r., it's out of the page.

b.) Determine the magnitude of the force.

$$\begin{aligned} F &= iLB \\ &= (30 \text{ A})(.12 \text{ m})(.9 \text{ T})\sin 60^\circ \\ &= 2.81 \text{ N} \end{aligned}$$

Example 7: (courtesy of Mr. White): A 10 cm 10 cm square loop of wire hangs vertically, as shown here. When the current in the wire is 0.245 A counterclockwise, a scale supporting the wire measures a downward force of 3.48×10^{-2} N. Find the magnitude of the magnetic field.



According to the *r-h-r*, the force will be downward. Its magnitude will be:

$$F = iLB$$

$$\Rightarrow 3.48 \times 10^{-2} \text{ N} = (.245 \text{ A})(.10 \text{ m})B \sin 90^\circ$$

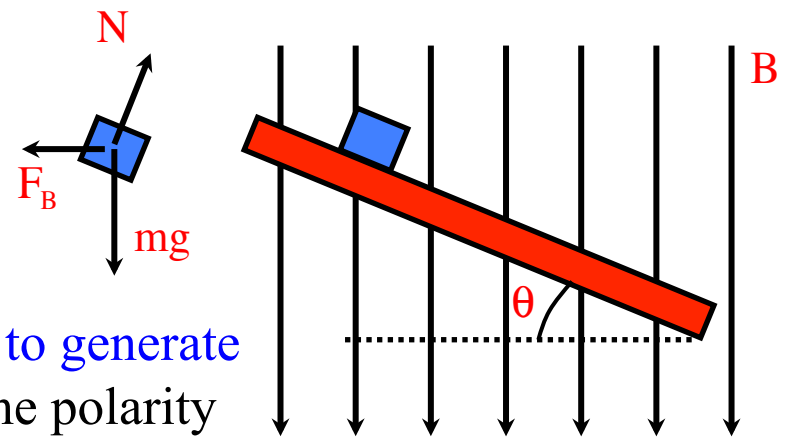
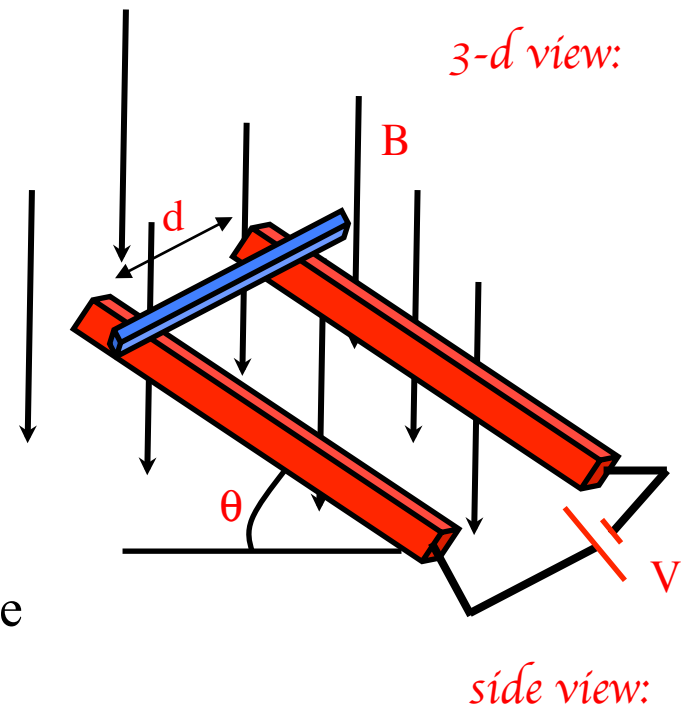
$$\Rightarrow B = 1.42 \text{ T}$$

Note: This could easily have been turned into a torque problem . . . just sayin . . .

Example 8: Two metal ramps (in red) at an angle θ are d units apart and are bathed in a downward B -fld. A battery wire is connected to each ramp making a circuit. A metal rod would slide frictionless down the ramp if it were not for the magnetic force provided by the B -fld. In fact, in this case the rod is motionless. Assuming the net resistance of the system is R :

a.) What must the polarity of the battery be if the rod is to stay stationary on the ramp (again, assume the rod is frictionless).

For equilibrium, you can see from observation that you need a component of the magnetic force *to the left* to counter the component of gravity *to the right*. Reverse engineering $\vec{F} = i\vec{L} \times \vec{B}$ yields the need for a current *into the page* to generate a magnetic force in that direction . . . So the polarity of the battery must be **high-side on left**.



b.) *How big must* the battery voltage be to effect this situation?

Doing a f.d.b. and breaking the forces into components, we can determine *i*:

$$\sum F_y :$$

$$-mg + N \cos \theta = 0$$

$$\Rightarrow N = \frac{mg}{\cos \theta}$$

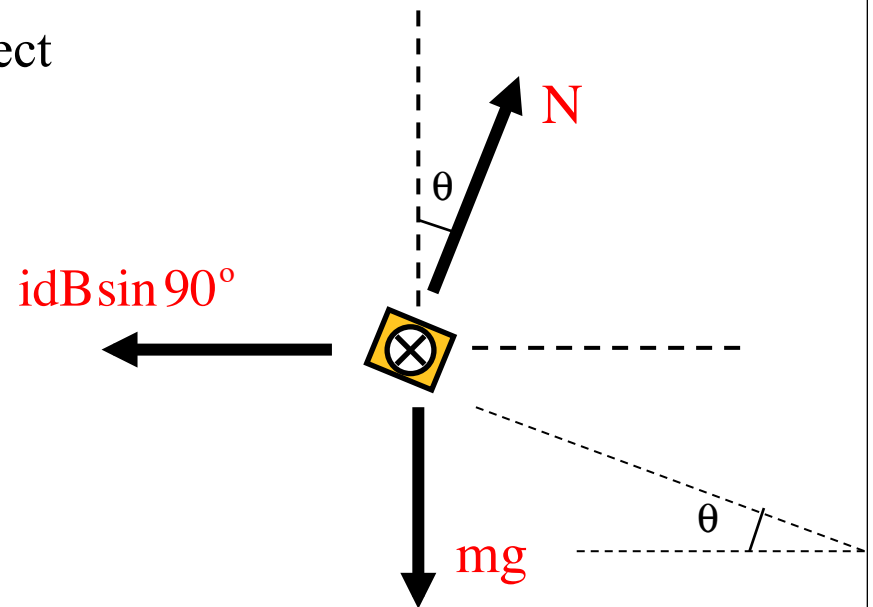
$$\sum F_x :$$

$$-i d B + N \sin \theta = 0$$

$$\Rightarrow i = \frac{N \sin \theta}{dB}$$

$$\Rightarrow i = \frac{\left(\frac{mg}{\cos \theta} \right) \sin \theta}{dB}$$

$$\Rightarrow i = \frac{mg \tan \theta}{dB}$$



From Ohm's Law:

$$V = iR$$

$$= \left(\frac{mg \tan \theta}{dB} \right) R$$

Another Point of Order

The relationship for the force on a wire in a B -fld looks like:

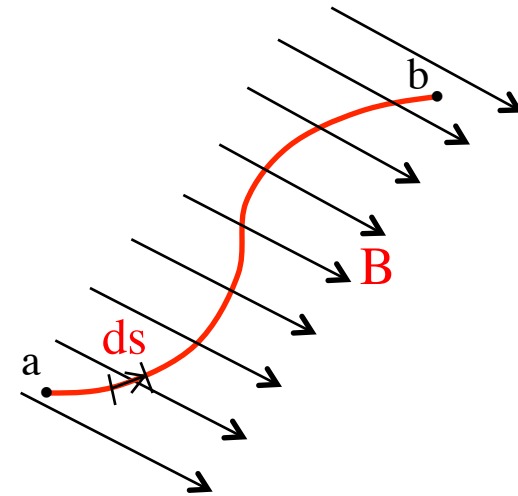
$$\vec{F}_{\text{wire}} = i\vec{L} \times \vec{B}$$

but this only works when L is a straight wire, B is constant and the angle between the two vectors remains the same throughout.

If any of those parameters vary, you need to use an integral. Specifically:

- define a differential displacement ds along the wire;
- cross ds into the magnetic field evaluated at ds ;
- execute the integral:

$$\vec{F}_{\text{wire}} = i \int_a^b d\vec{s} \times \vec{B}$$



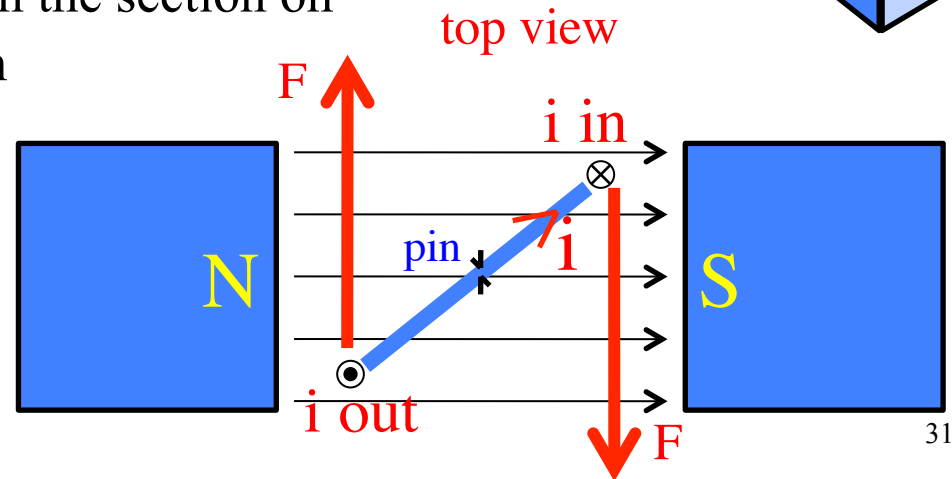
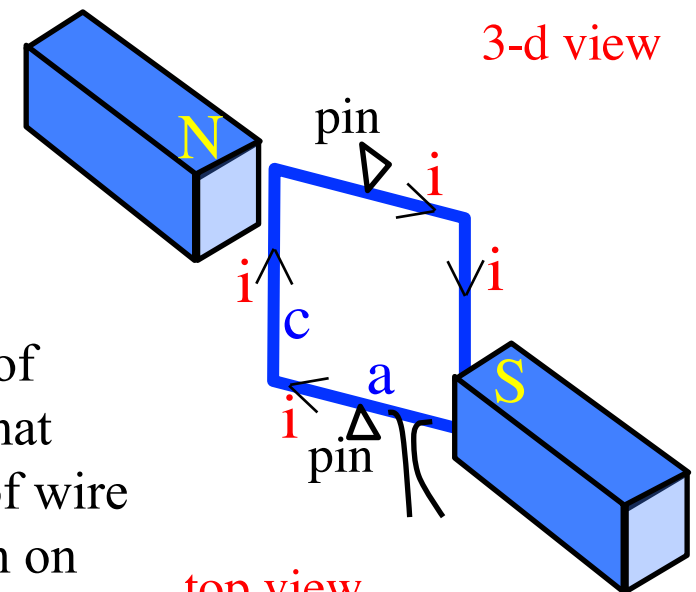
Galvanometer and Torque on a Current-Carrying Loop

Consider a current-carrying loop of width a and height c , pinned at the top and bottom, bathed in a magnetic field.

How will the current i respond to the B -fld?

Look down from the top and ignore the billowing of magnetic field generated by the bar magnets. In that view, the current moves across the upper section of wire as shown, moves down into the page in the section on the right, and moves out of the page in the section on the lower left.

The *r.h.r.* suggests the direction of the magnetic forces on those sections are as shown.

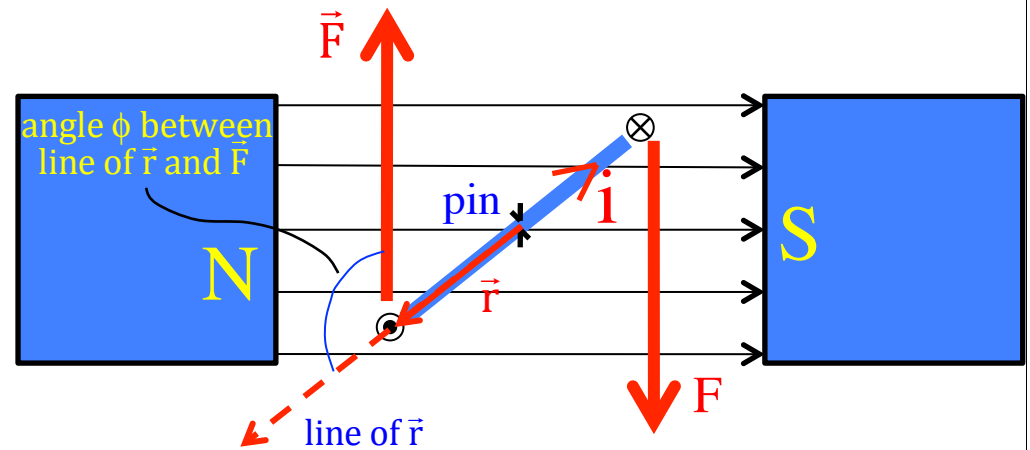


The *magnitude* of the *magnetic force* will be:

$$\begin{aligned}\vec{F} &= i\vec{L} \times \vec{B} \\ \Rightarrow F &= iLB \sin 90^\circ \\ &= icB\end{aligned}$$

The *total torque* due to the current in the wires into and out of the page is:

$$\begin{aligned}|\vec{\tau}| &= 2|\vec{r} \times \vec{F}| \\ &= 2|\vec{r}||\vec{F}|\sin\phi \\ &= 2\left(\frac{a}{2}\right)(icB)\sin\phi \\ &= i(ac)B\sin\phi\end{aligned}$$



where ϕ is the angle between the *line of the force* \vec{F} and the *line of the position vector* \vec{r} .

Two observations:

1.) There could be N loops in our coil, so the most general torque expression should be the torque expression we've already derived multiplied by N .

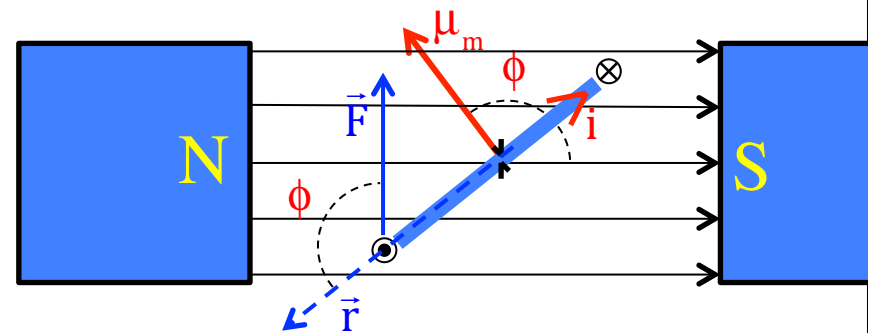
2.) Notice that ac is the area A of the loop.

With this, we can write the torque calculation as:

$$\begin{aligned} |\vec{\tau}| &= N [i(ac) B \sin \phi] \\ &= N [iAB \sin \phi] \end{aligned}$$

--Although this is not something I think the AP folks are likely to ever test on, consider:

Lay your right-hand on the loop in the direction of the current. The direction your outstretched thumb points will define a direction perpendicular to the face of the coil.



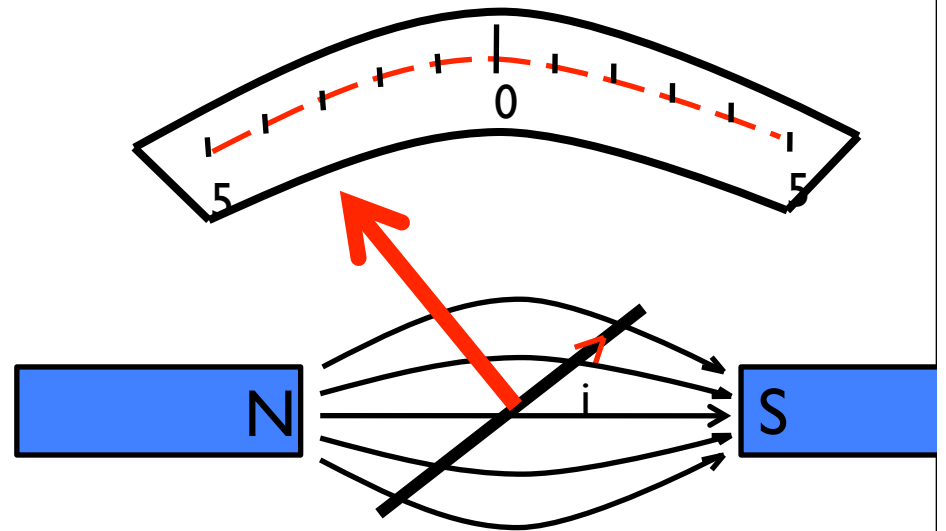
Define a vector, called the *magnetic moment* $\vec{\mu}_M$, whose direction is in that perpendicular direction and whose magnitude is equal to NiA . Crossing that magnetic moment vector into the B-fld yields a magnitude:

$$|\vec{\mu} \times \vec{B}| = (NiA) B \sin \phi \quad \text{which is the torque on the coil} \dots$$

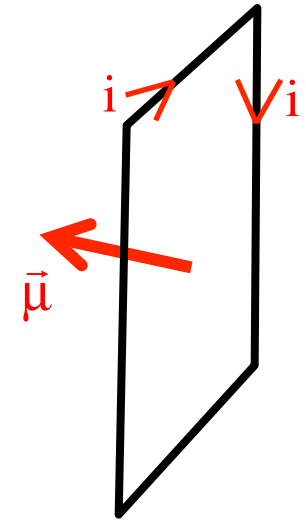
Aside from giving physics teachers an excuse for having students do torque calculations in an E&M section, the real usefulness of all of this quite cool.

If you put a **pinned coil** in a **magnetic field**, **attach a spring** to it to provide a **restoring torque**, then **attach a pointer** hung over a **scale**, you have the **makings of a meter**.

And, in fact, that is exactly **how a GALVANOMETER is made**—it's a coil **suspended in a magnetic field** with a **spring attached to it to provide a restoring torque**, so that when you put 5×10^{-4} **amps** through it, the **torque provided** by the moving current in the *B-fl*d coupled with the restoring torque sees the **pointer swing “full deflection” over the scale** . . . and that is very cool . . .



Example 8: (courtesy of Mr. White): A rectangular coil of dimensions 5.40 cm x 8.50 cm consists of 25 turns of wire and carries a current of 15.0 mA.



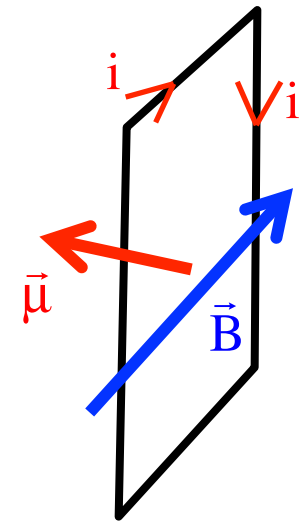
a.) Calculate the magnetic moment of the coil, putting the direction onto the sketch.

$$\begin{aligned} |\vec{\mu}| &= NiA \\ &= (25)(15 \times 10^{-3} \text{ A})[(.054 \text{ m})(.080 \text{ m})] \\ &= 1.72 \times 10^{-3} \text{ A} \cdot \text{m}^2 \end{aligned}$$

Circling the right hand in the direction of the current yields a thumb-point leftward (the direction should be perpendicular to the face of the coil)

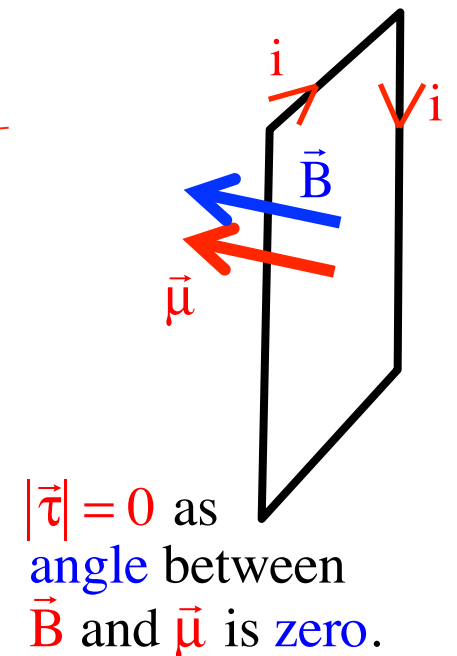
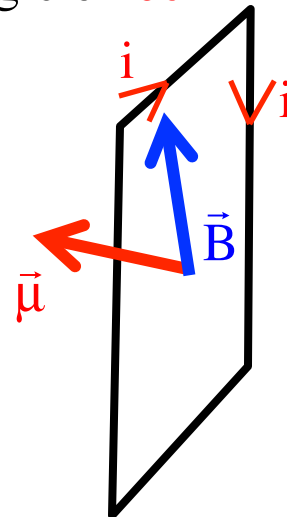
b.) If a magnetic field of 0.350 T is applied parallel to the plane of the loop, what is the magnitude of the torque acting on the loop?

$$\begin{aligned}
 |\vec{\tau}| &= |\vec{\mu} \times \vec{B}| \\
 &= |\vec{\mu}| |\vec{B}| \sin \theta \\
 &= (1.72 \times 10^{-3} \text{ T} \cdot \text{m}^2)(0.35 \text{ T}) \sin 90^\circ \\
 &= 6.02 \times 10^{-4} \text{ N} \cdot \text{m} \quad (\text{direction downward})
 \end{aligned}$$



c.) Calculate the magnitude of the torque on the coil when the 0.350 T B-field makes an angle of 60° with $\vec{\mu}$.

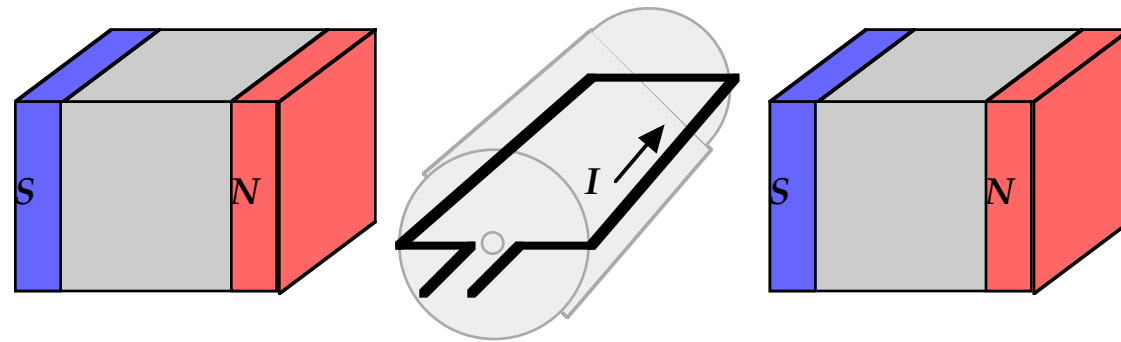
$$\begin{aligned}
 |\vec{\tau}| &= |\vec{\mu}| |\vec{B}| \sin \theta \\
 &= (1.72 \times 10^{-3} \text{ T} \cdot \text{m}^2)(0.35 \text{ T}) \sin 60^\circ \\
 &= 5.21 \times 10^{-4} \text{ N} \cdot \text{m}
 \end{aligned}$$



$|\vec{\tau}| = 0$ as angle between \vec{B} and $\vec{\mu}$ is zero.

Electric Motors (courtesy of Mr. White)

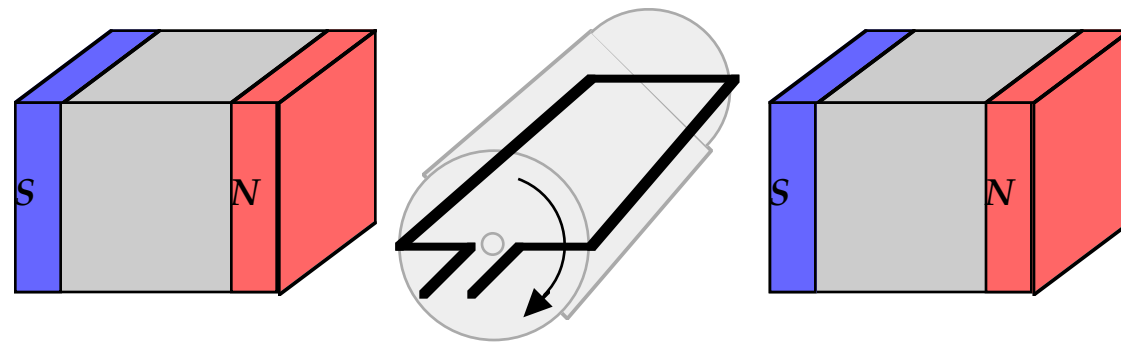
Electric motors convert electrical energy to kinetic energy, and are created by placing a current-carrying loop in an external magnetic field. There are a number of different ways of doing this, but here is one common type:



Note from Fletch: So “current in” motivates rotation (hence a motor).

Electric Generators (courtesy of Mr. White)

Electric generators convert mechanical energy (work provided by an external source) to electrical energy. Motors and generators, in most cases, have the same physical structure. (Note from Fletch: In other words, as far as energy conversion goes, a *motor* is just a *generator* run “in reverse.”)



Note from Fletch: So rotation motivates a current (hence a generator).

Other Devices

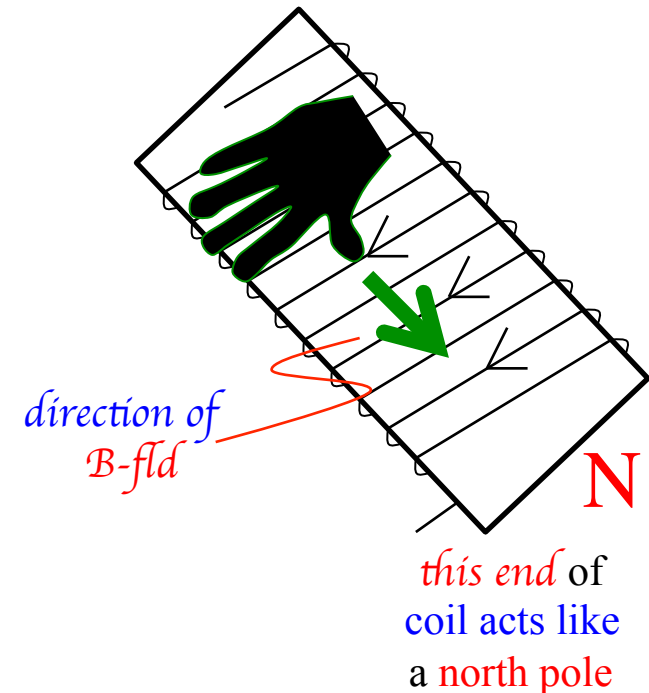
A *little more* sophisticated version of a motor required one bit of information that would normally not be covered until next chapter.

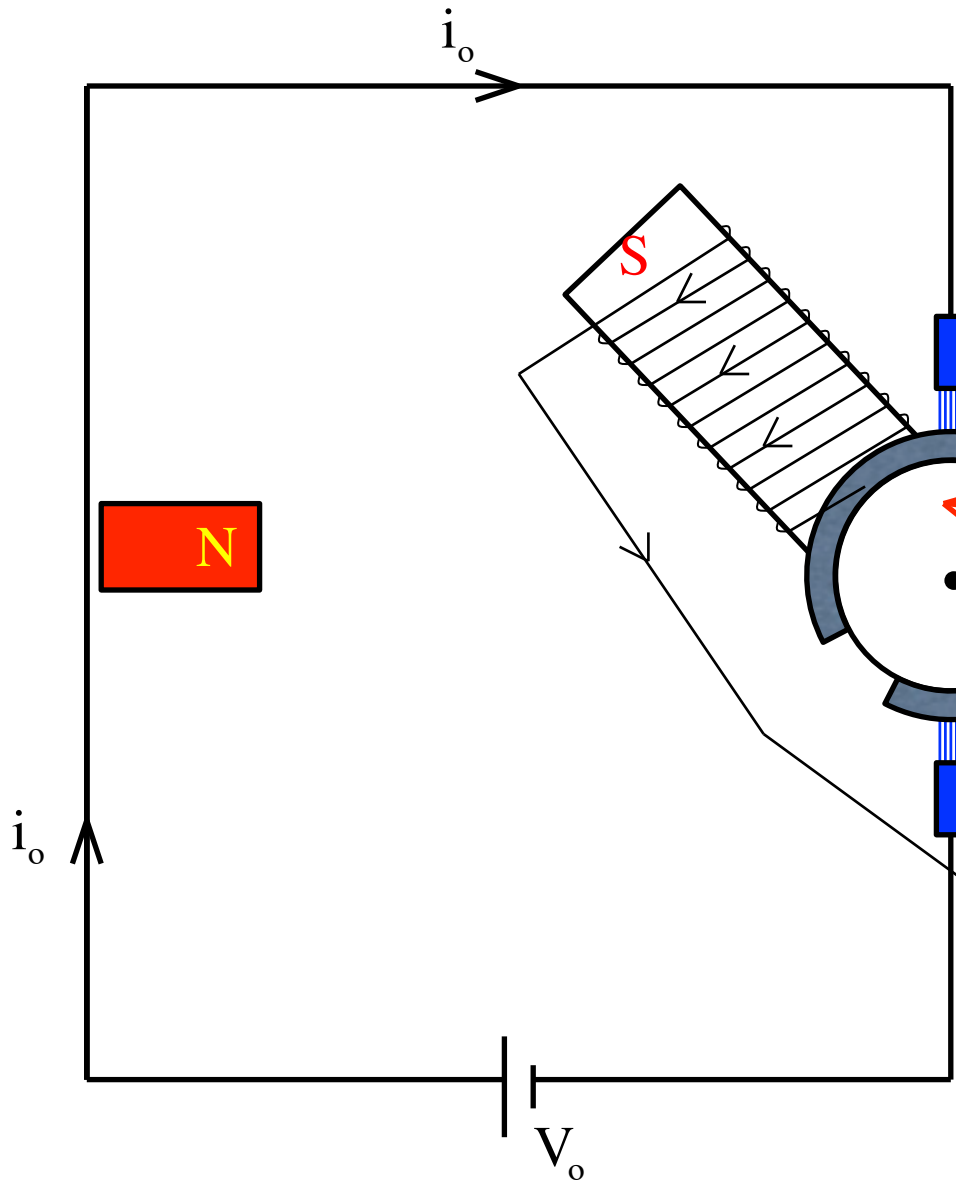
It is charge in motion that generates magnetic fields. With current carrying coils, the generated magnetic fields are *down the axis* and *through the face of the coil*.

A *handy trick* to determine the direction of a current carrying wire's *B-fl* is to lay your **right hand** on the coil with your fingers **following the direction of current in the coil**. The direction in which your extended right-thumb points identifies the direction of the coil's *B-fl*.

Note that with the *B-fl* extending along the axis as it does, the **coil's ends look like north and south poles**.

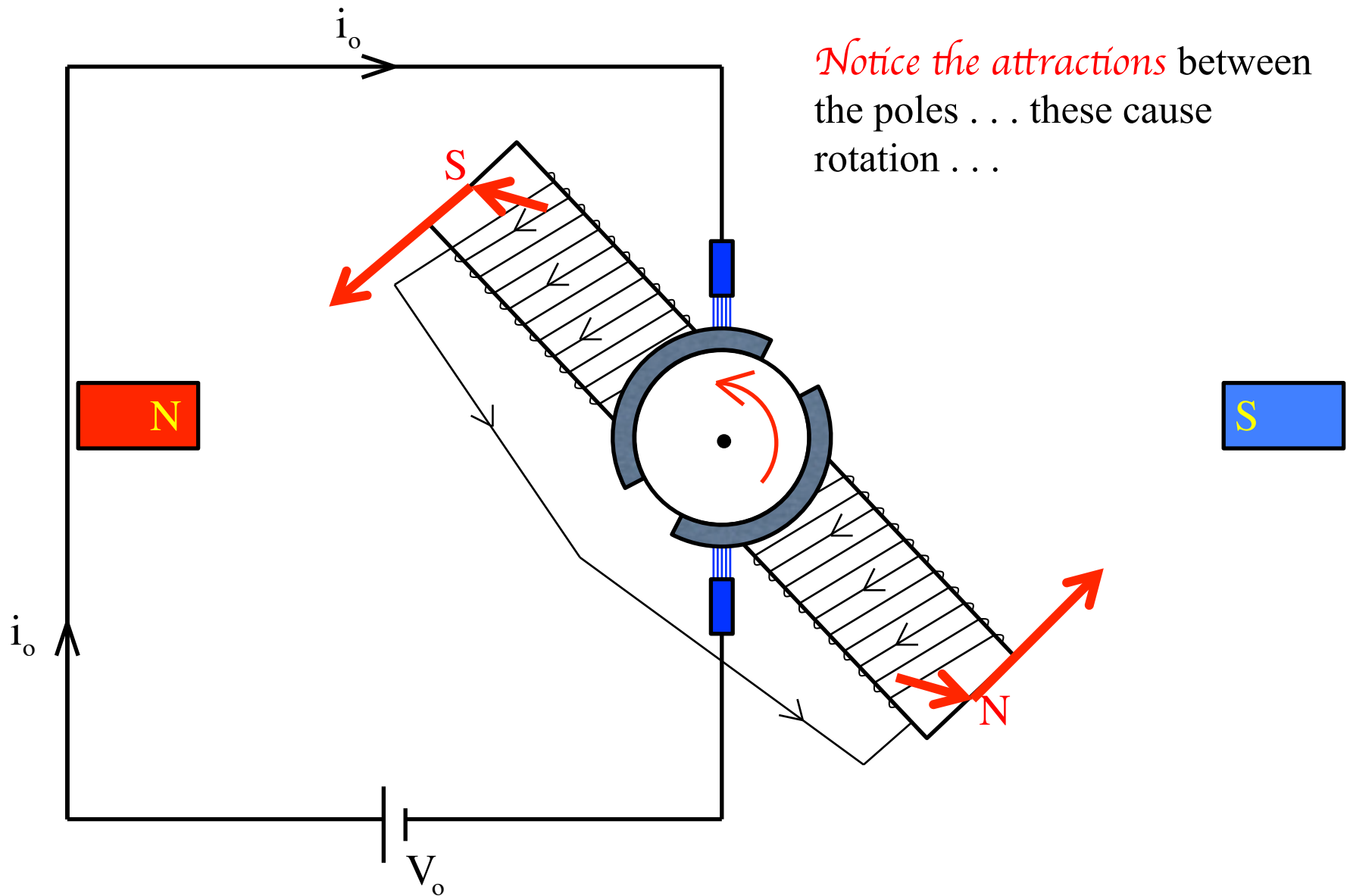
With that, consider the following:



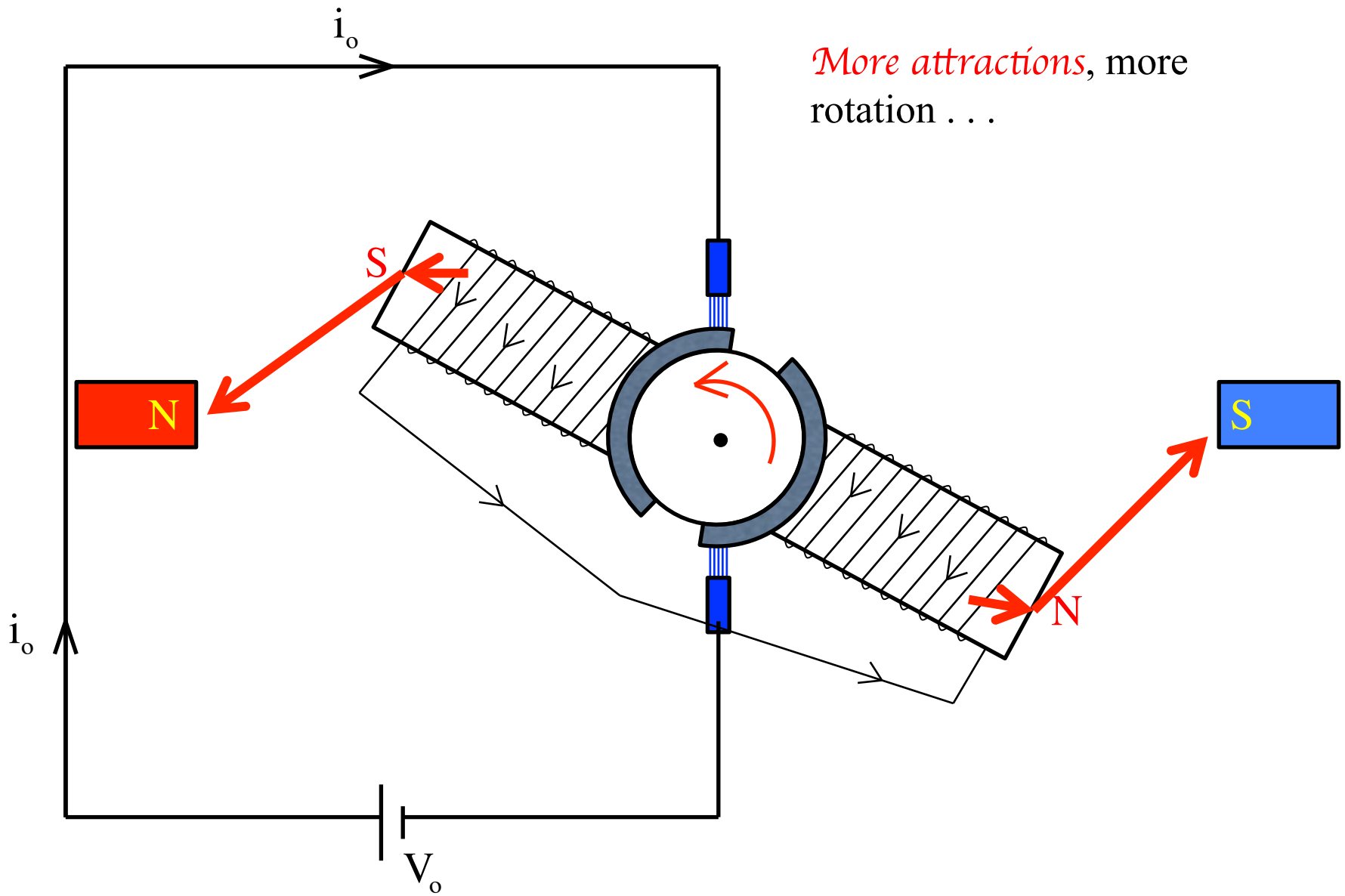


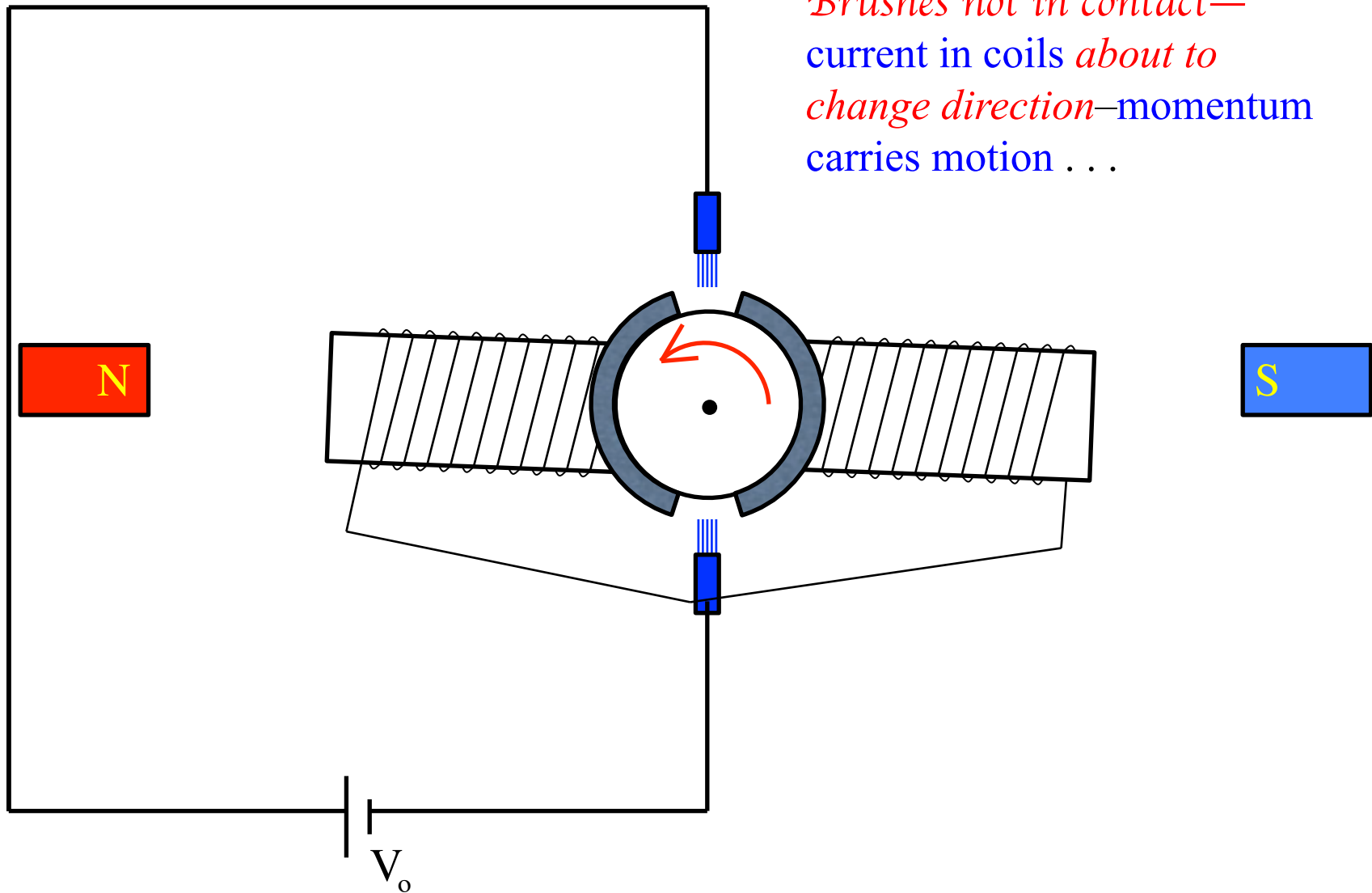
Follow the current
 from the battery, through the
 brushes to the coil, then
 determine the *B-flid* due to the
 coil's current (see below).

Fingers of right hand **curl**
 along line of current;
 thumb identifies
 direction of B-
 field (with B-
 field lines **exiting**
 this end, so it's a **North**
Pole)

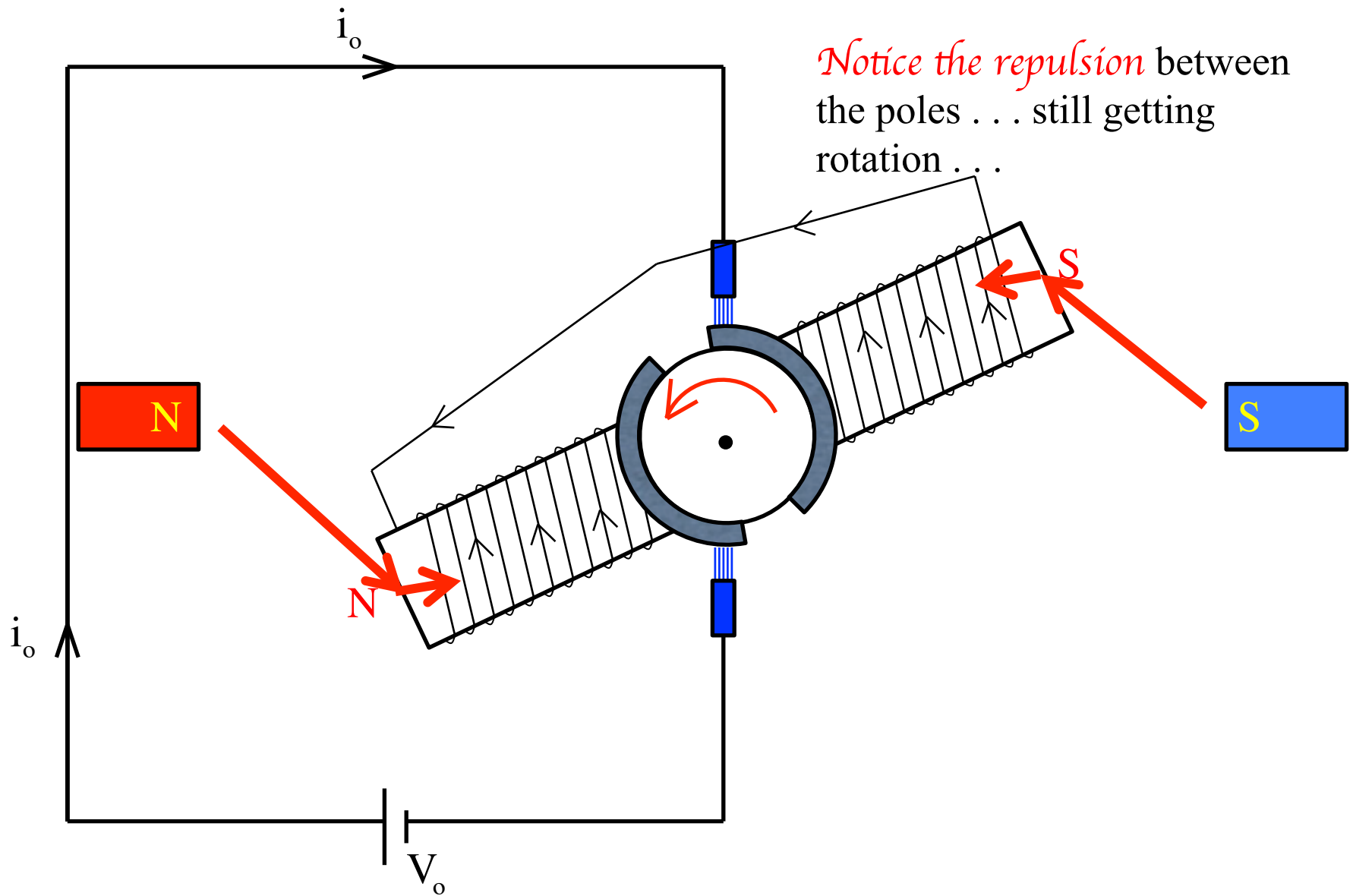


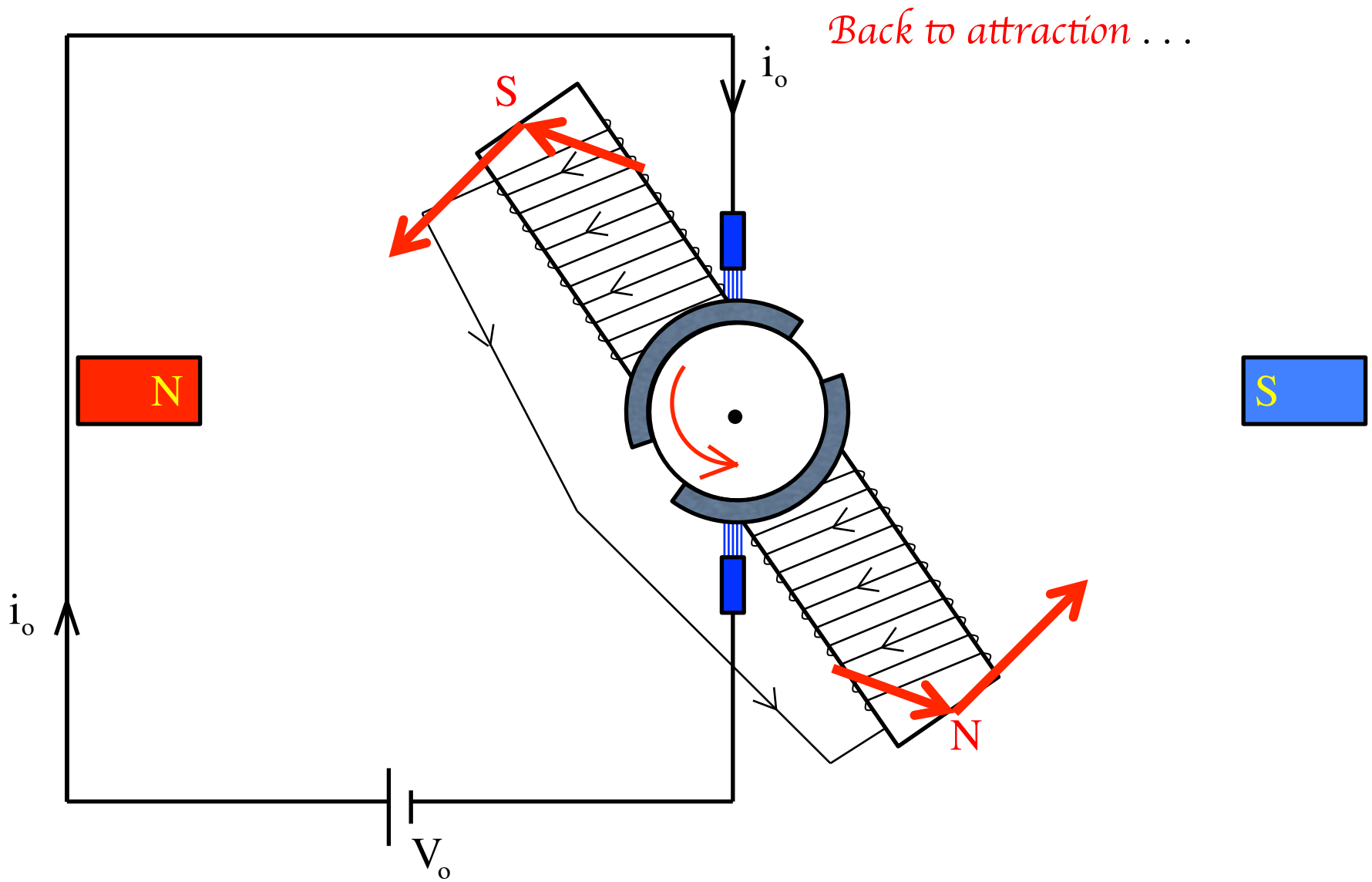
Notice the attractions between the poles . . . these cause rotation . . .





*Brushes not in contact—
current in coils about to
change direction—momentum
carries motion . . .*



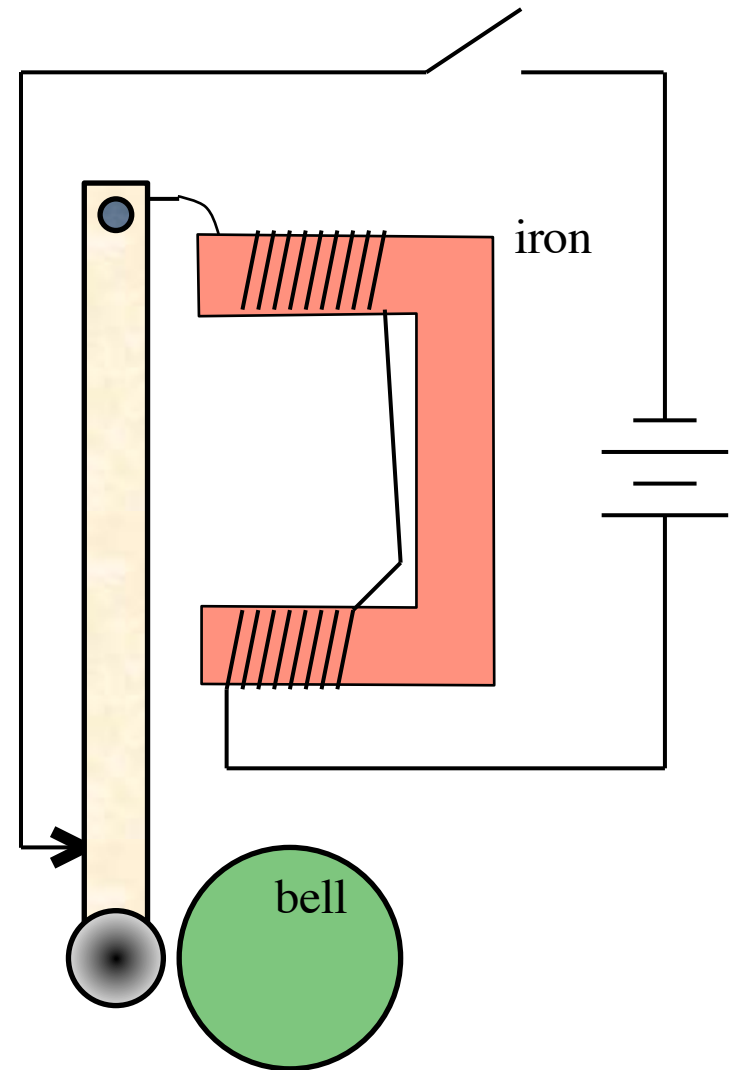


Back to attraction ...

This is a DC motor...

What is it?

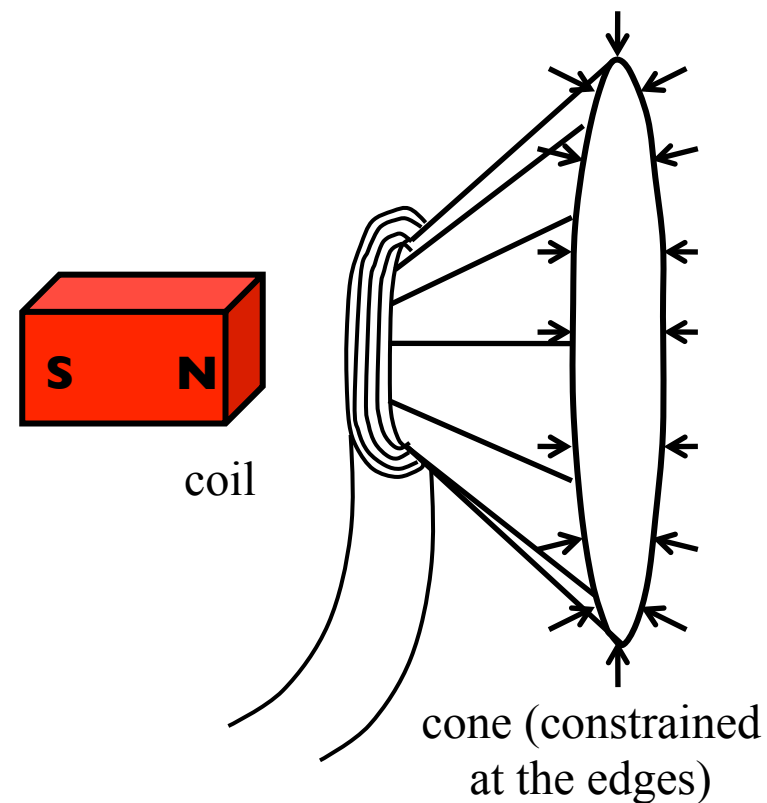
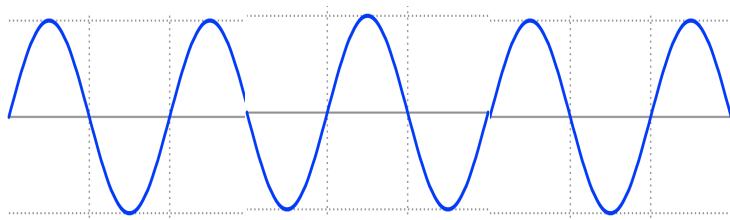
What you are looking at here is the **circuit** for an **old-fashioned door bell**. See if you can follow through to see how the mechanism works (**note the direction of the magnetic field in the green horseshoe electromagnet**, and the direction of the induced magnetic field in the blue bar opposite the poles of the horseshoe magnetic, when the current flows).

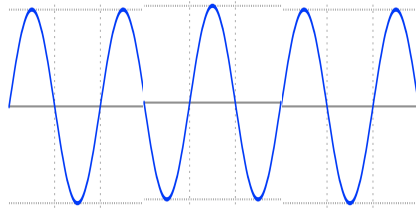


What is it?

This is the design for a **loud speaker** (actually, the coil and magnet are usually switched in real line). This is how it works:

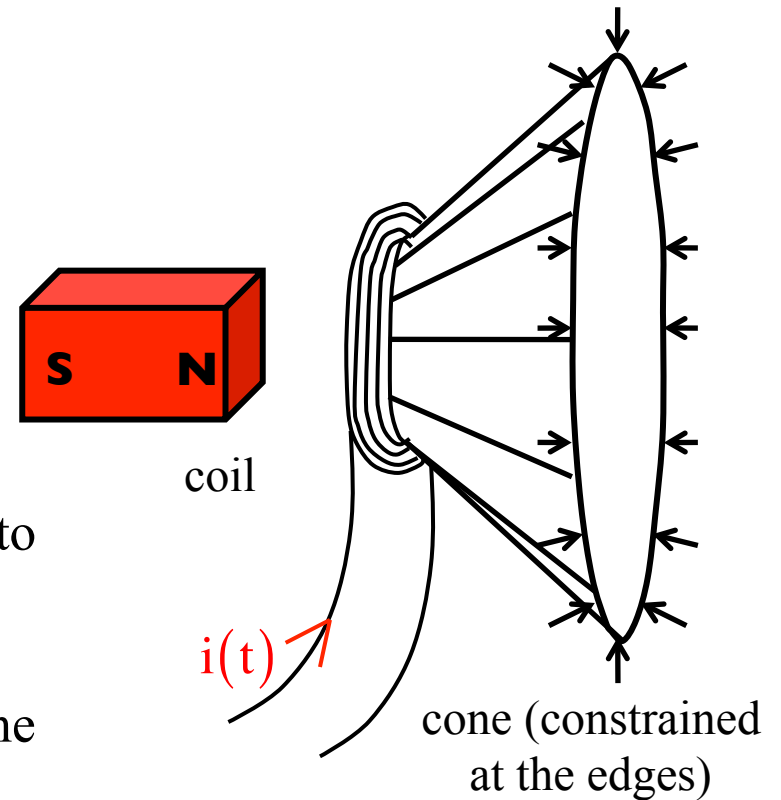
a.) Let's assume we want to project a **256 Hz (middle C) sound wave** into the room. The signal would look like the sine wave shown below.





b.) During the first half of the cycle, let's assume the direction of the time-varying current through the coil is directed as shown on the sketch. There are two things to note:

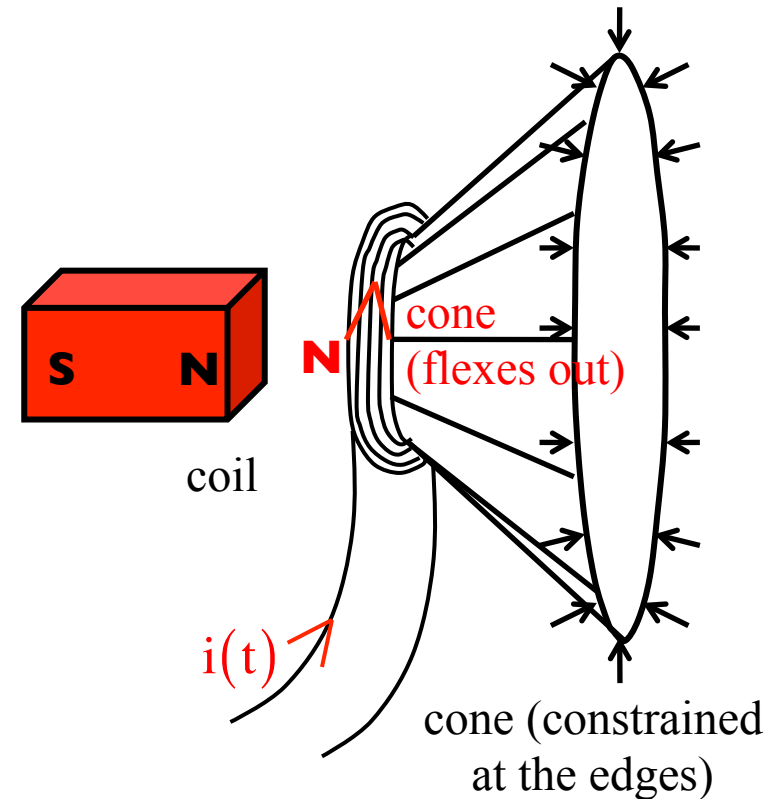
i.) Being sinusoidal, the current will increase to a maximum, then will drop back down to zero whereupon the direction will change and the current will again increase to a maximum in the opposite direction, then proceed back to zero. This pattern will continue through time.



ii.) With the current moving in the direction noted, the direction of the induced magnetic field in the coil (alternate right-hand rule) will leave the side of the coil closest to the permanent magnet a **North Pole** (see sketch).

c.) The north pole of the permanent magnet will interact with the induced north pole of the current carrying coil, and the net effect will be a **repulsion** experienced by both the coil AND the cone. As the cone is fixed at its outside edge, this will flex the cone outward with the amount of flex being dependent upon the size of the current at the given instant.

d.) As the cone flexes outward, it will compress air into a high pressure region. That pressure ridge will travel away from the cone at approximately **330 m/s**, or the speed of sound in air.

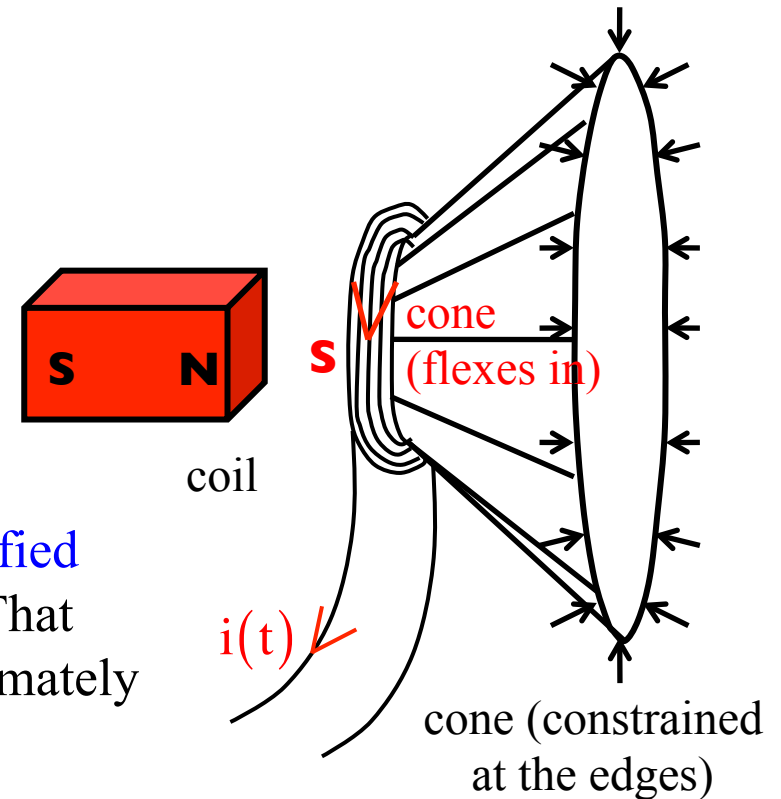


e.) As the current proceeds down toward zero, the cone will relax, pulling back.

f.) When the current direction changes, the direction of the induced magnetic field in the coil will change and the “pulling back” will proceed through equilibrium and into a flexing inward. The degree of flexing will, again, depend upon how much current is moving through the coil at a given instant.

g.) As the cone flexes inward, it will create a rarified region of air generating a low pressure region. That region will travel away from the cone at approximately 330 m/s, or the speed of sound in air.

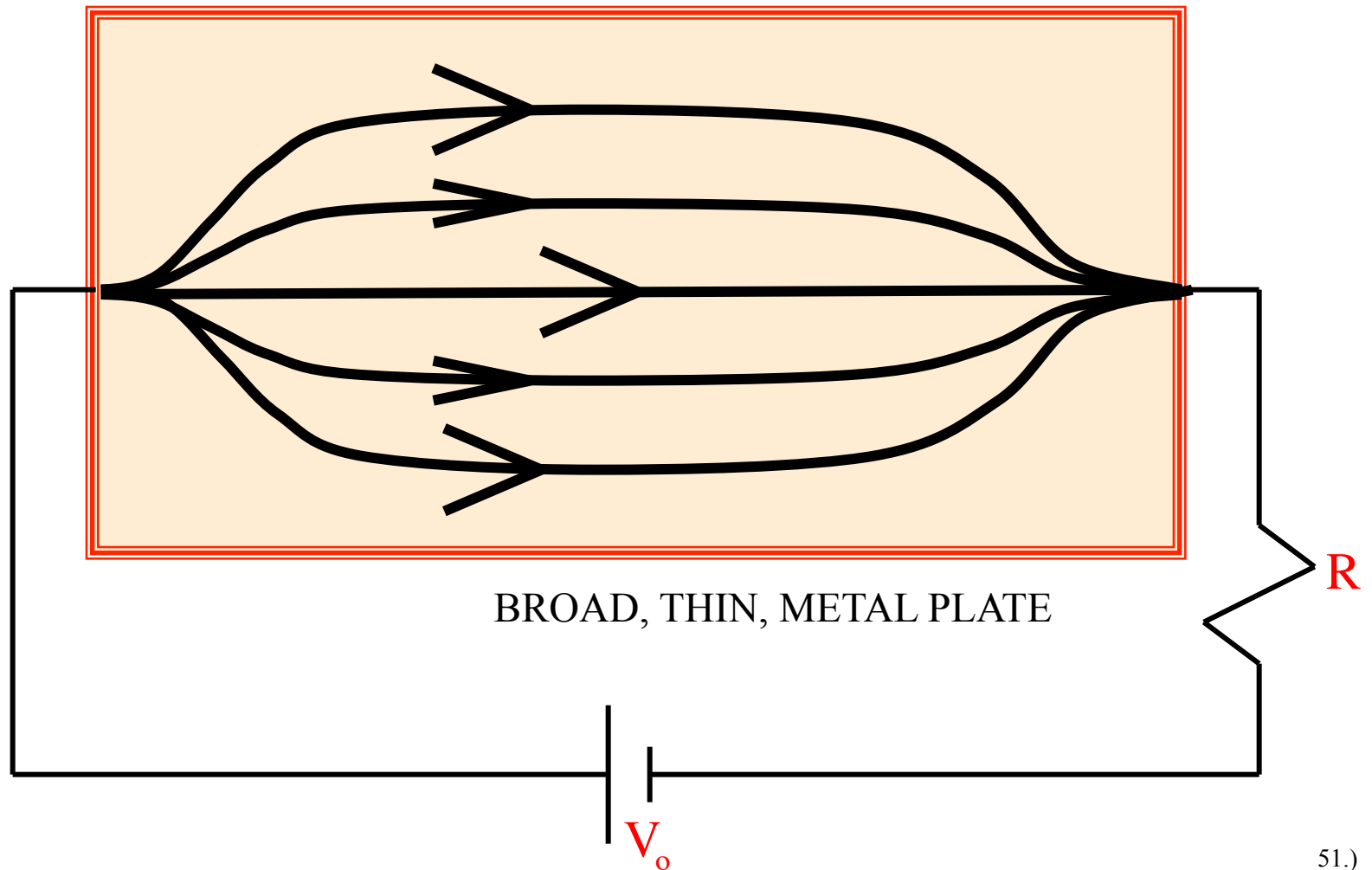
h.) This flexing outward, then inward, then outward will occur at the current frequency, or 256 Hz in our example, and the pressure variations will pass by your ear at a frequency of 256 Hz. That, in turn, will wiggle the little hairs in your ears creating electrical impulses that your brain will interpret as sound. Clever of nature, eh?



HALL EFFECT

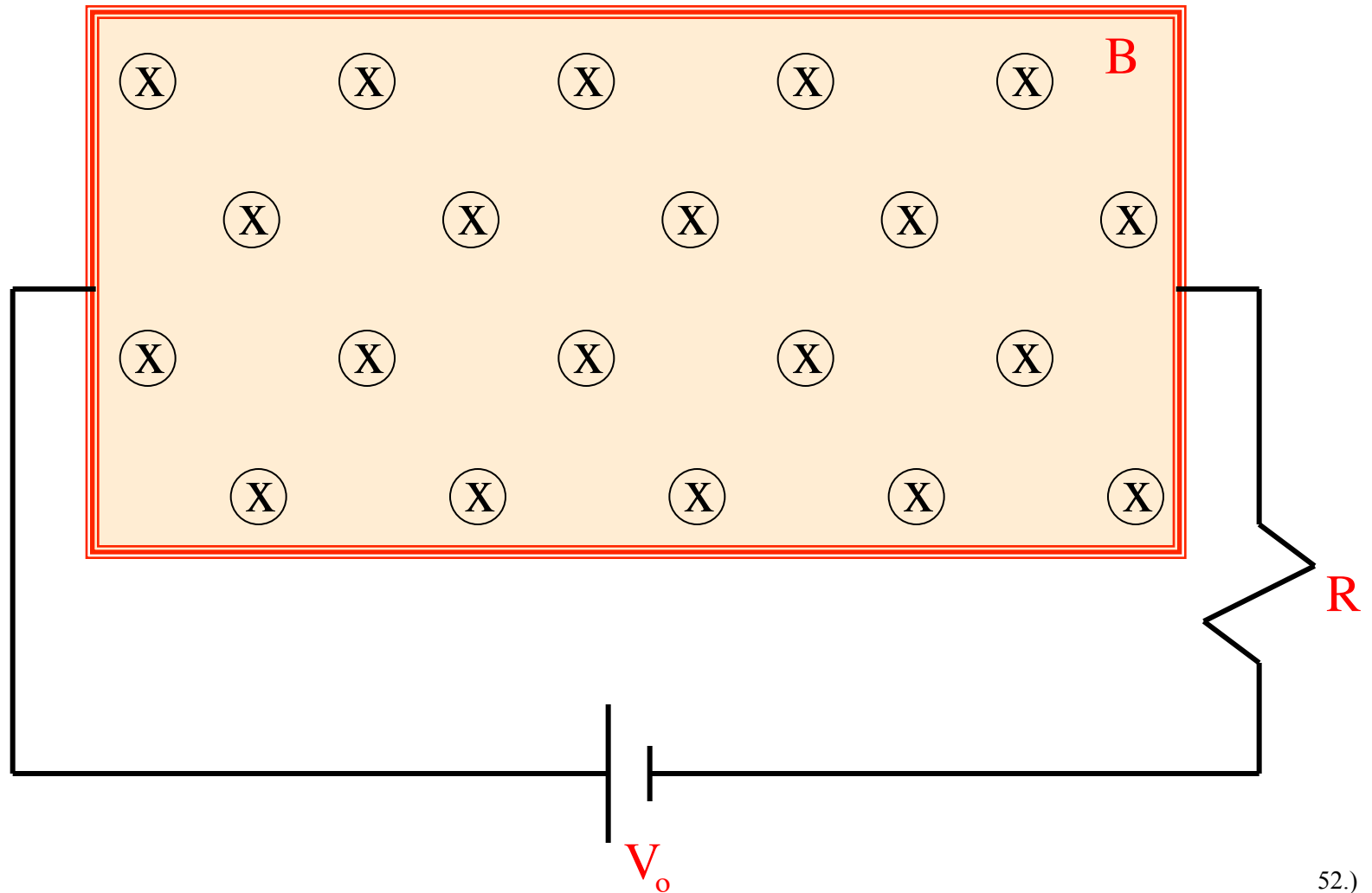
An experiment to prove that *negative charges* moves through electrical circuits when current flows.

How normal current flows through the plate.



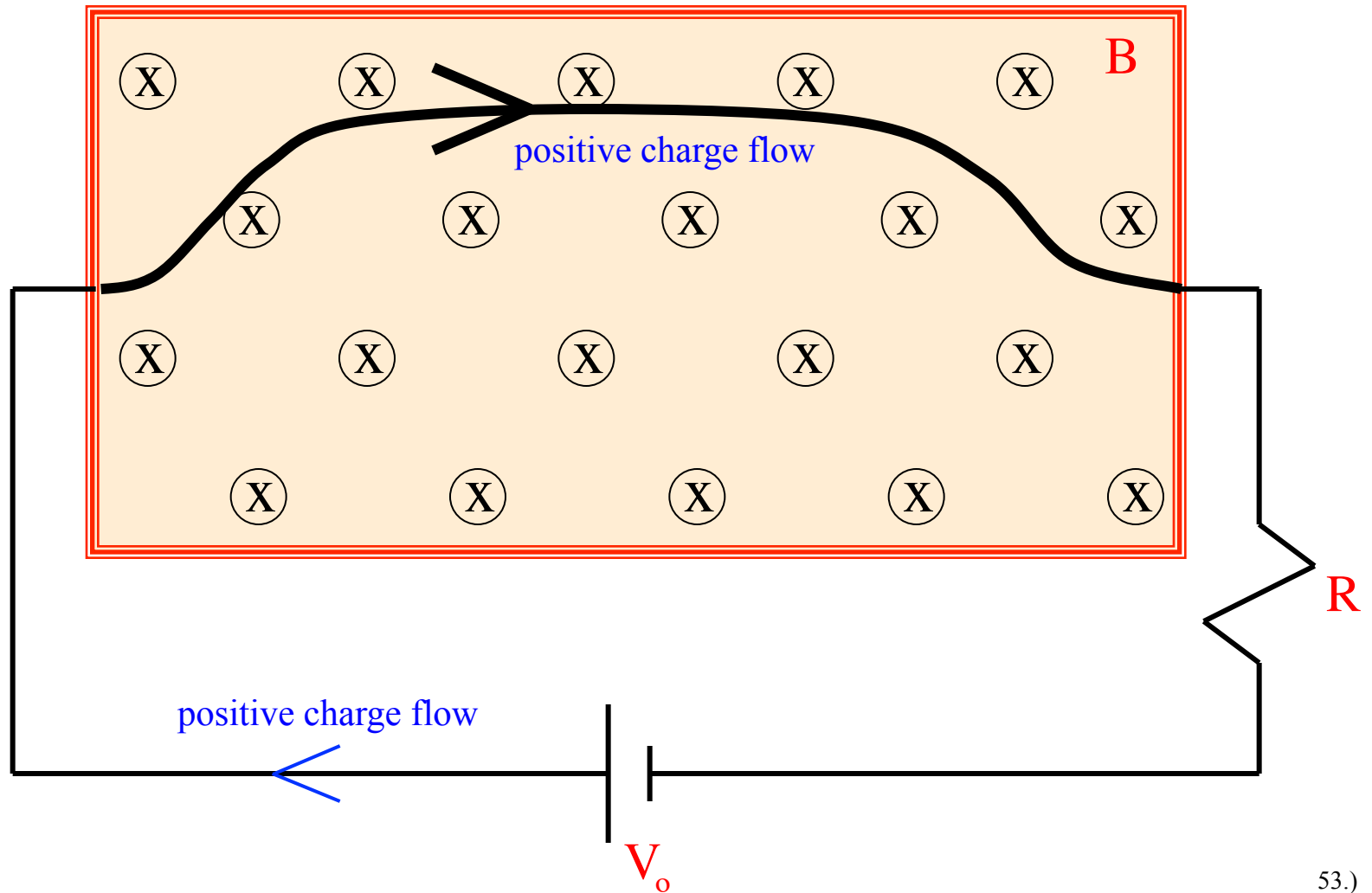
What happens when the plate is permeated by *B*-fld.

*plate
permeated
by B-fld*



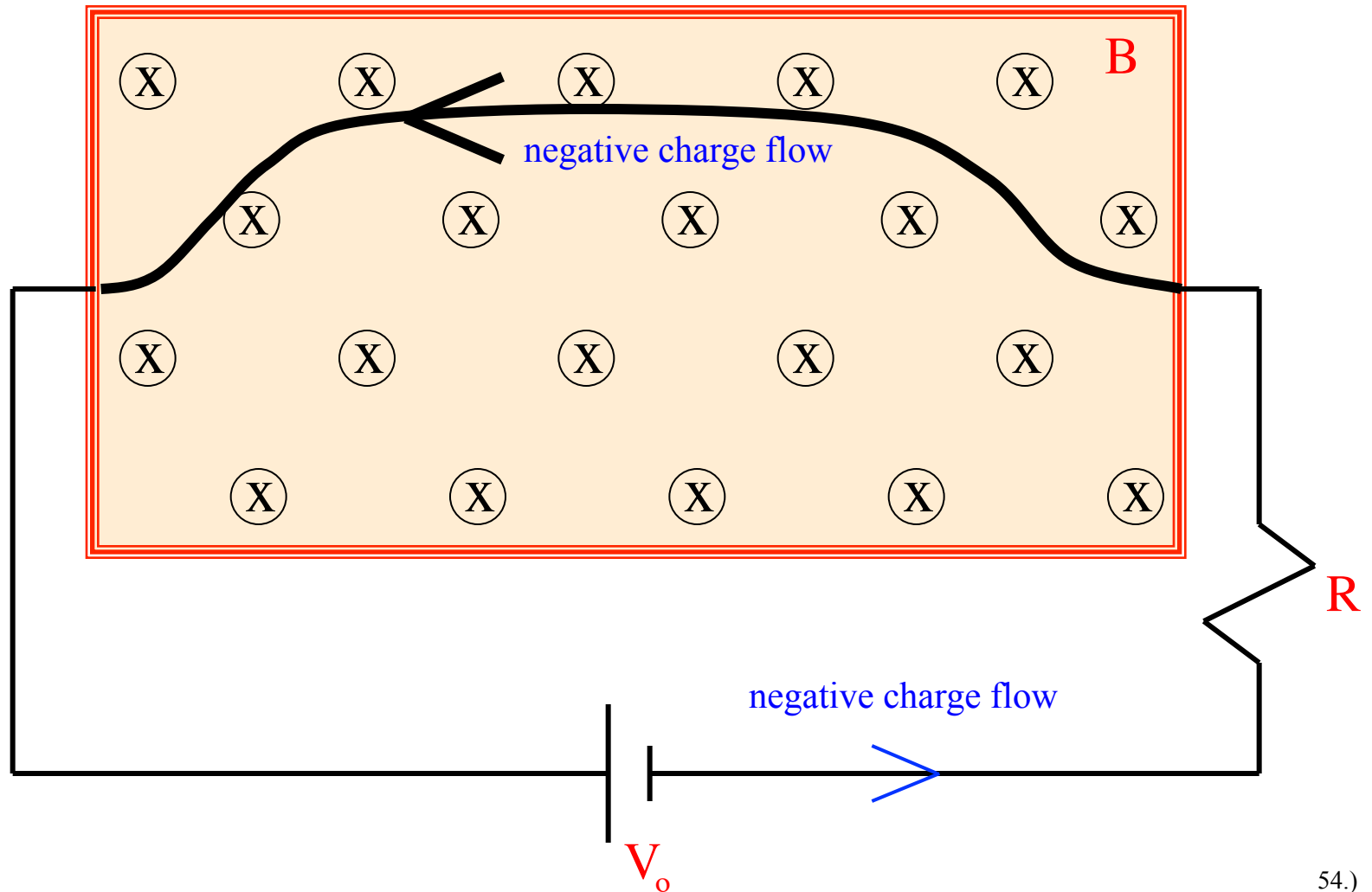
... if *POSITIVE* charge is assumed to flow through circuit?

positive charge flow suggests upper side of plate will be *high voltage side*



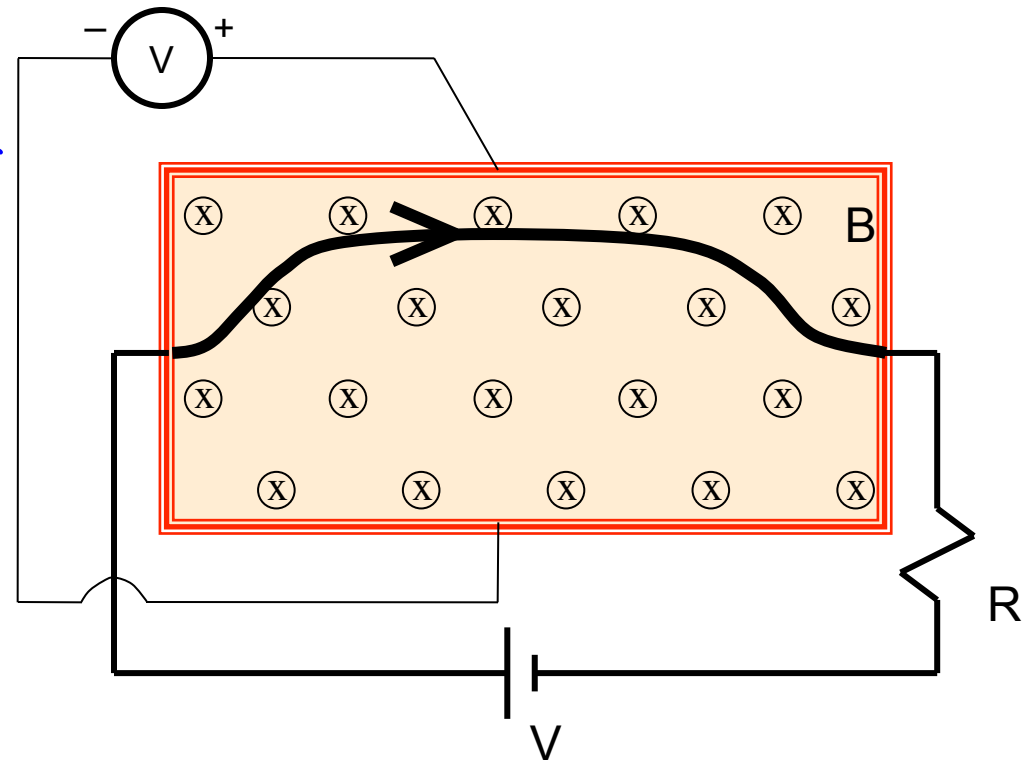
... if *NEGATIVE* charge is assumed to flow through circuit?

negative charge flow suggests upper side of plate will be *low voltage* side



So back to the positive charge-flow assumption:

If there exists a **predominance of positive charge** on the upper side of the plate making that side of the plate **higher voltage** than the bottom side, then **placing a voltmeter with its + and - terminals as positioned** would produce a **meter reading that was sensible** (that is, the needle would swing in the appropriate direction).



In fact, **if this experiment was done**, the **needle would swing in the wrong direction**. In other words, **the meter is hooked up wrong**. What are accumulating **on the upper side of the plate** are **not positive charges**, they are **NEGATIVE charges**.

In short, it is **negative charge** that flows through circuits.

3 Types of Magnetism

Ferromagnetism - material has a permanent “magnetic moment,” due to microscopic “domains” in which moments are aligned.

Paramagnetism - materials has a small magnetic susceptibility, that only becomes evident when placed in an external magnetic field.

Diamagnetism - material does not have permanent magnetic moments. In the presence of an external magnetic field, a weak magnetic moment is induced in a direction opposite to the external field.